

### **X-force**

$$X = - \oint_l \mathbf{p} \cdot d\mathbf{y} = - \oint_l (\mathbf{p} - \mathbf{p}_\infty) \cdot d\mathbf{y} = - \oint_l \left( \frac{\rho_\infty U_\infty^2}{2} - \frac{\rho_\infty U^2}{2} \right) \cdot d\mathbf{y} = \frac{\rho_\infty}{2} \oint_l U^2 \cdot d\mathbf{y}$$

### **Y-force**

$$Y = \oint_l \mathbf{p} \cdot d\mathbf{x} = \oint_l (\mathbf{p} - \mathbf{p}_\infty) \cdot d\mathbf{x} = \oint_l \left( \frac{\rho_\infty U_\infty^2}{2} - \frac{\rho_\infty U^2}{2} \right) \cdot d\mathbf{x} = -\frac{\rho_\infty}{2} \oint_l U^2 \cdot d\mathbf{x}$$

$$X - iY = - \oint_l \mathbf{p} \cdot d\mathbf{y} + i p dx = \dots = i \frac{\rho_\infty}{2} \oint_l U^2 \cdot d\bar{z}$$

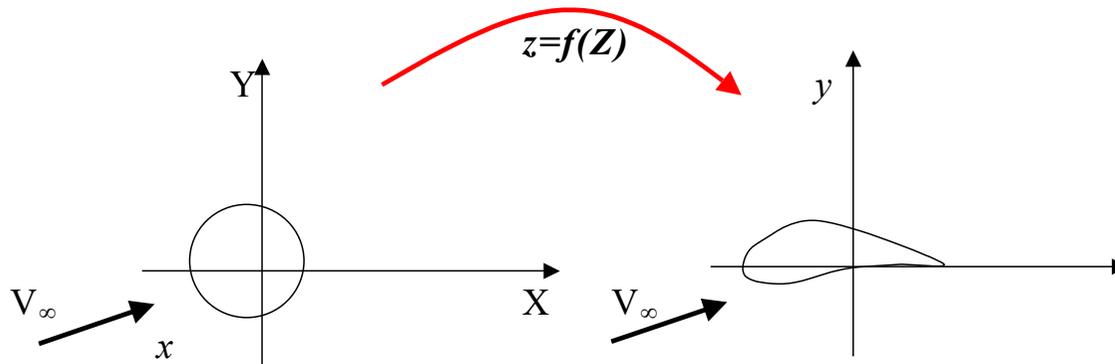
$$M_0 = - \oint_l \mathbf{p} \cdot (y d\mathbf{y} + x d\mathbf{x}) = - \oint_l \mathbf{p} \cdot \text{Re}(z d\bar{z}) = \dots = \frac{\rho_\infty}{2} \text{Re} \oint_l U^2 \cdot z d\bar{z}$$

$$\frac{dw}{dz} = |U|(\cos \Theta - i \sin \Theta) = |U| \frac{d\bar{z}}{ds}$$

$$\left(\frac{dw}{dz}\right)^2 = \left(|U| \frac{d\bar{z}}{ds}\right)^2 = U^2 \frac{d\bar{z} \cdot d\bar{z}}{ds \cdot ds} = U^2 \frac{d\bar{z} \cdot \boxed{d\bar{z}}}{dz \cdot \boxed{d\bar{z}}} = U^2 \frac{d\bar{z}}{dz}$$

$$U^2 d\bar{z} = \left(\frac{dw}{dz}\right)^2 dz \quad ; \quad U^2 z d\bar{z} = \left(\frac{dw}{dz}\right)^2 z dz$$

$$X - iY = i \frac{\rho_\infty}{2} \oint_l U^2 \cdot d\bar{z} = i \frac{\rho_\infty}{2} \oint_l \left(\frac{dw}{dz}\right)^2 dz \quad ; \quad M_0 = \frac{\rho_\infty}{2} \operatorname{Re} \oint_l \left(\frac{dw}{dz}\right)^2 z dz$$



$$X - iY = i \frac{\rho_\infty}{2} \oint_l \left( \frac{dw}{dz} \right)^2 dz = i \frac{\rho_\infty}{2} \oint_l \left( \frac{dw}{dZ} \right)^2 \frac{dZ}{dz} \frac{dZ}{dz} dz = i \frac{\rho_\infty}{2} \oint_l \left( \frac{dw}{dZ} \right)^2 \frac{dZ}{dz} dZ$$

$$M_0 = \frac{\rho_\infty}{2} \operatorname{Re} \oint_l \left( \frac{dw}{dz} \right)^2 z dz = \frac{\rho_\infty}{2} \operatorname{Re} \oint_l \left( \frac{dw}{dZ} \right)^2 \frac{dZ}{dz} \frac{dZ}{dz} z dz = \frac{\rho_\infty}{2} \operatorname{Re} \oint_l \left( \frac{dw}{dZ} \right)^2 \frac{dZ}{dz} z dZ$$

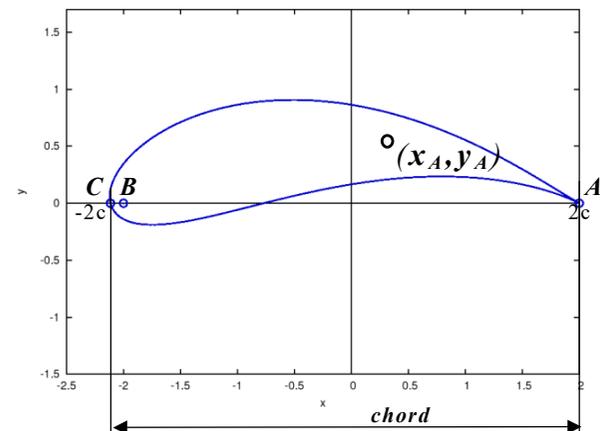
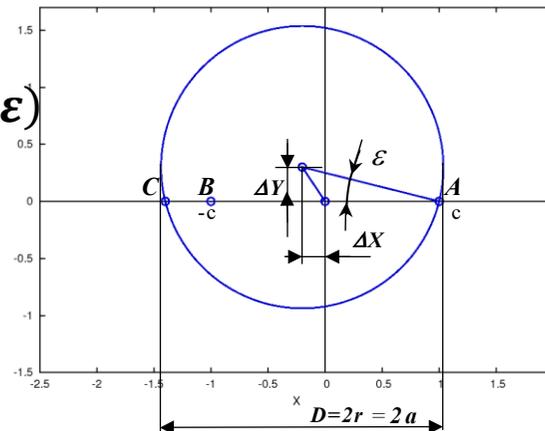
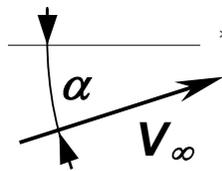
$$M_A = M_0 + Yx_A - Xy_A = M_0 + \operatorname{Re} [i(X - iY)(x_A + iy_A)]$$

$$z = f(Z) = Z + c_0 + c_1 Z^{-1} + c_2 Z^{-2} + c_3 Z^{-3} + \dots$$

$$\frac{dZ}{dz} = 1 + c_1 Z^{-2} + \dots \quad ; \quad \frac{dZ}{dz} z = Z + c_0 + 2c_1 Z^{-1} + \dots$$

$$W(Z) = U_\infty \left( e^{-i\alpha} Z + \frac{e^{i\alpha}}{Z} \right) + i \frac{\Gamma}{2\pi} \ln(Z)$$

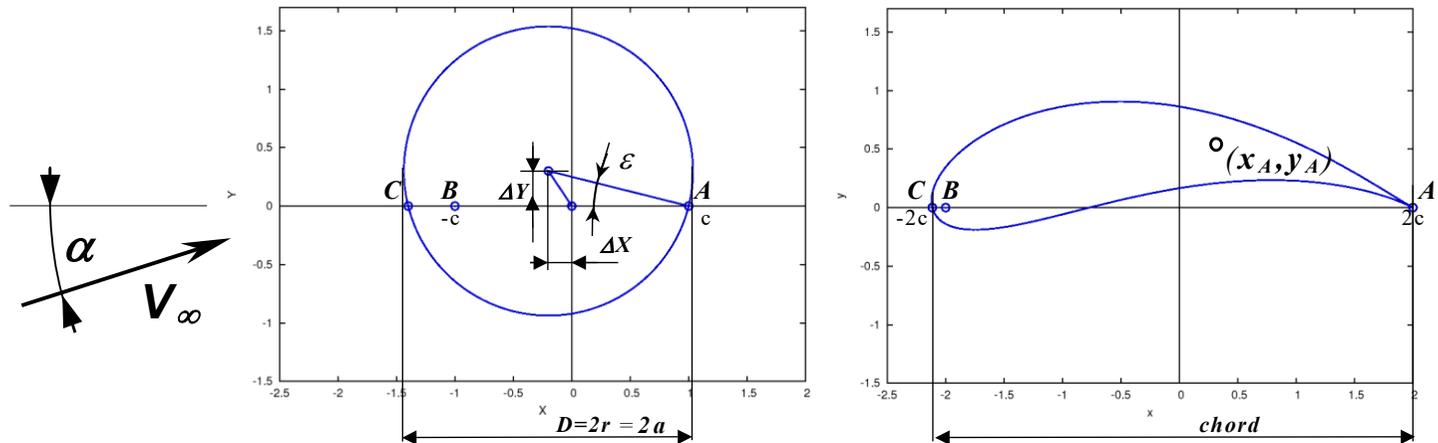
$$\Gamma = 4\pi a U_\infty \cdot \sin(\alpha - \varepsilon)$$



$$M_A = \frac{\rho_\infty}{2} \operatorname{Re} \oint_l \left( \frac{dw}{dZ} \right)^2 \frac{dZ}{dz} (z - z_A) dZ = \dots$$

$$= 2\pi \rho_\infty U_\infty^2 \operatorname{Re} \{ i e^{-2i\alpha} [c_1 + a(z_A + c_0) e^{-i\varepsilon}] - i a(z_A + c_0) e^{i\varepsilon} \}$$

$$z_A = c_0 - \frac{c_1}{a} e^{i\varepsilon} \Rightarrow M_A = \text{const}$$

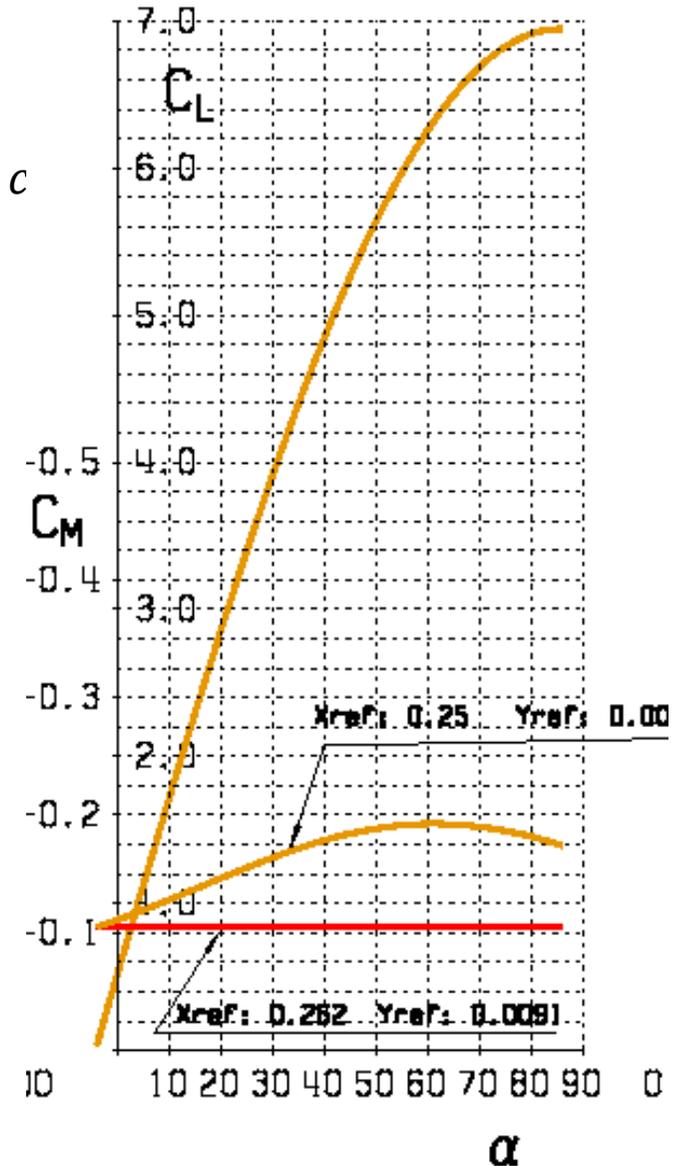


# AERODYNAMIC CENTER

...  $M_A = c$

thin, symmetrical airfoil  $x_{ac}/c = 0.25$   
 $y_{ac}/c = 0.00$

thick, symmetrical/  
 nonsymmetrical airfoil  $x_{ac}/c \approx 0.25$   
 $y_{ac}/c \approx 0.00$



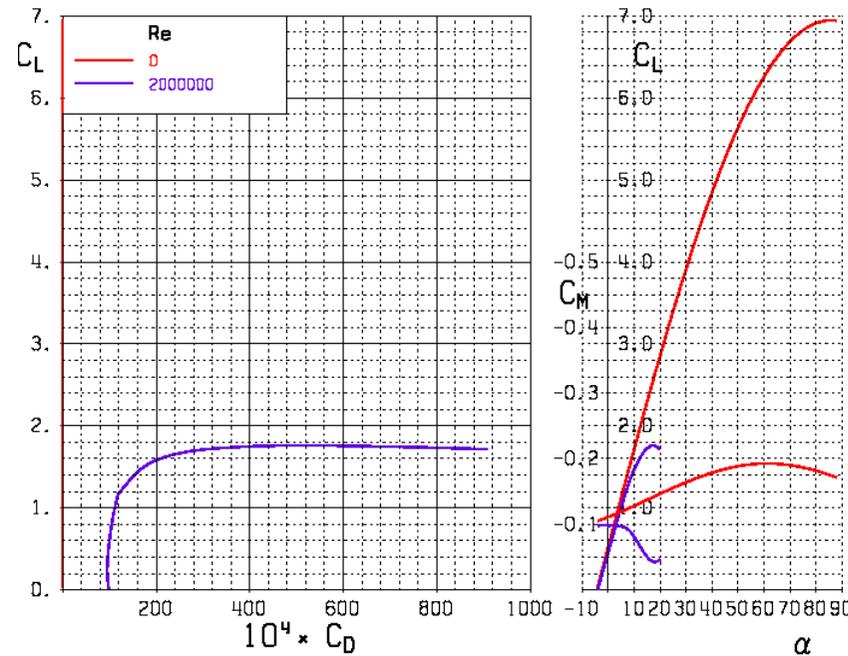
airfoil: NACA 4412

potential flow

$$L = C_L \cdot S \cdot \rho_\infty \frac{V_\infty^2}{2} \quad C_L = \text{fun}(\alpha, \text{geom}, Re, \dots)$$

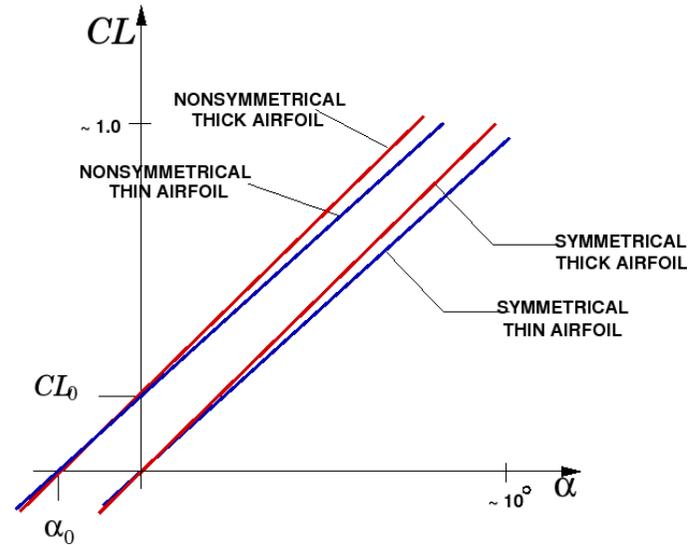
$$D = C_D \cdot S \cdot \rho_\infty \frac{V_\infty^2}{2} \quad C_D = \text{fun}(\alpha, \text{geom}, Re, \dots) [\text{FRICTION}]$$

$$My = Cm \cdot c \cdot S \cdot \rho_\infty \frac{V_\infty^2}{2} \quad Cm = \text{fun}(\alpha, \text{geom}, Re, (x_r, y_r) \dots)$$



$C_L, C_D, C_m \dots$

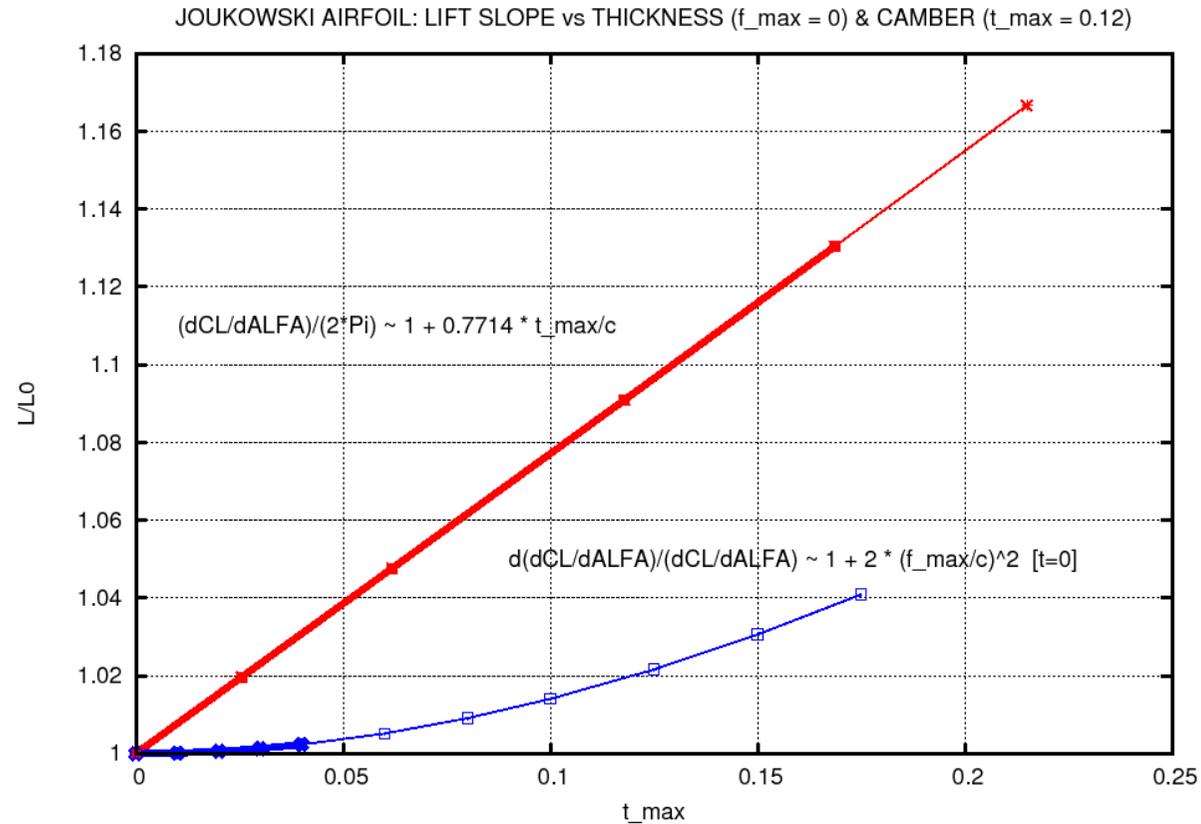
## low angles of attack



$$C_L \approx \frac{dC_L}{d\alpha} \cdot (\alpha - \alpha_0) = \frac{dC_L}{d\alpha} \cdot \alpha + C_{L0}$$

$$\frac{dC_L}{d\alpha} = \text{fun}(\bar{t}, \bar{f}, \text{Re}, \dots) \approx 0.1 \left[ \frac{1}{\text{deg}} \right]; 2\pi \left[ \frac{1}{\text{rad}} \right]$$

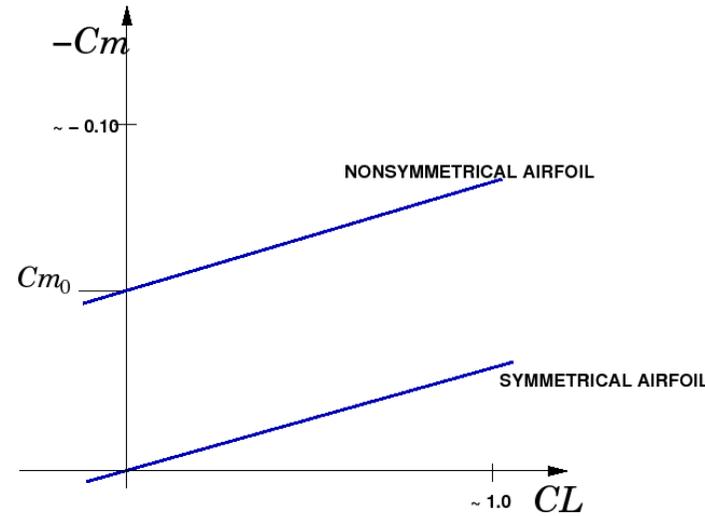
## lift force slope, Joukowski airfoil, inviscid flow



$$\left(\frac{dC_L}{d\alpha}\right)_{\bar{t}} - \left(\frac{dC_L}{d\alpha}\right)_{\bar{t}=0} = \text{fun1}(\bar{t})$$

$$\left(\frac{dC_L}{d\alpha}\right)_{\bar{f}} - \left(\frac{dC_L}{d\alpha}\right)_{\bar{f}=0} = \text{fun2}(\bar{f})$$

## low angles of attack



$$Cm \simeq Cm_0 + \frac{dCm}{dC_L} \cdot C_L$$

$$\begin{aligned} Cm_{(x_r, y_r)} &= Cm_{(x_0, y_0)} + C_Y \cdot (x_r - x_0)/c - C_X \cdot (y_r - y_0)/c \simeq \\ &\simeq Cm_{(x_0, y_0)} + C_L \cdot (x_r - x_0)/c \end{aligned}$$

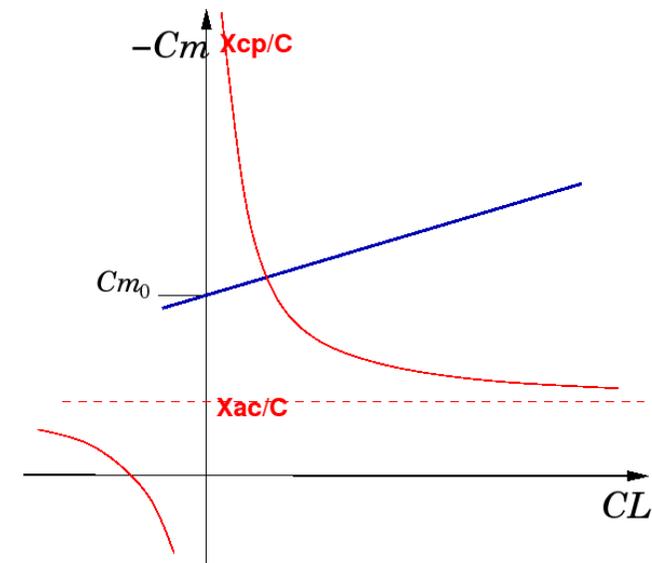
$$\left( \frac{dCm}{dC_L} \right)_{(x_r, y_r)} \simeq \left( \frac{dCm}{dC_L} \right)_{(x_0, y_0)} + (x_r - x_0)/c$$

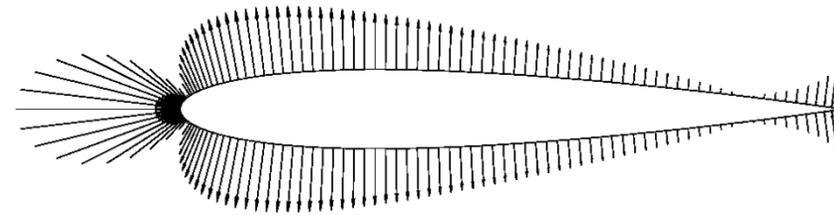
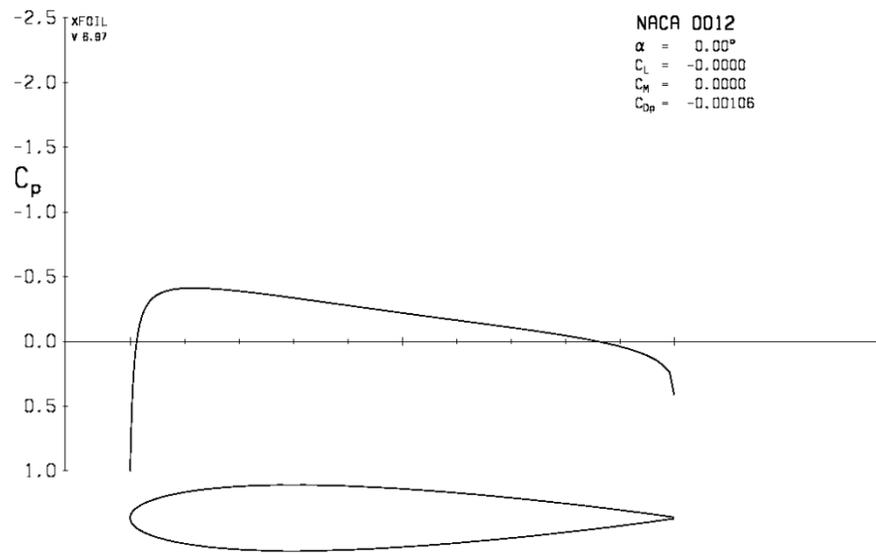
$$\left( \frac{dCm}{dC_L} \right)_{(x_{sa}, y_{sa})} = 0 \Rightarrow (x_{ac}/c) = (x_0/c) - \left( \frac{dCm}{dC_L} \right)_{(x_0, y_0)}$$

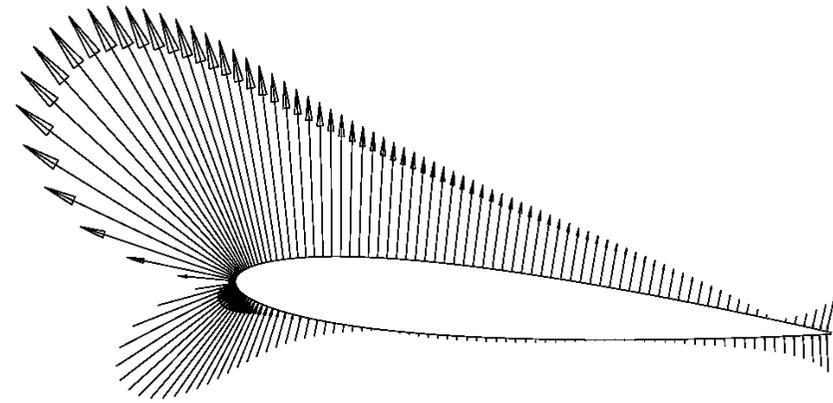
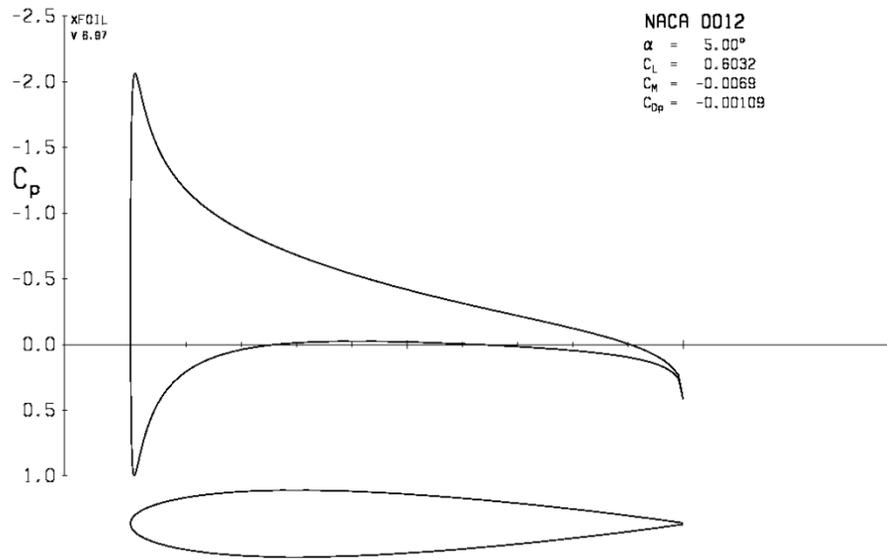
$$\begin{aligned}
 C_{m_{(x_r, y_r)}} &= C_{m_{(x_0, y_0)}} + C_Y \cdot (x_r - x_0)/c - C_X \cdot (y_r - y_0)/c \approx \\
 &\approx C_{m_{(x_0, y_0)}} + C_L \cdot (x_r - x_0)/c
 \end{aligned}$$

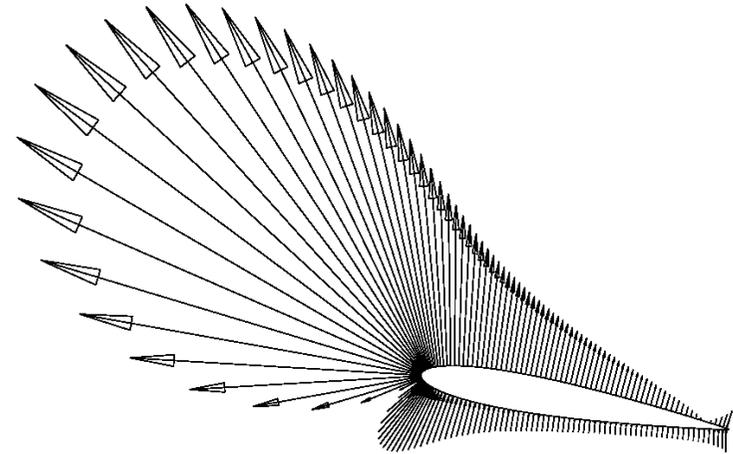
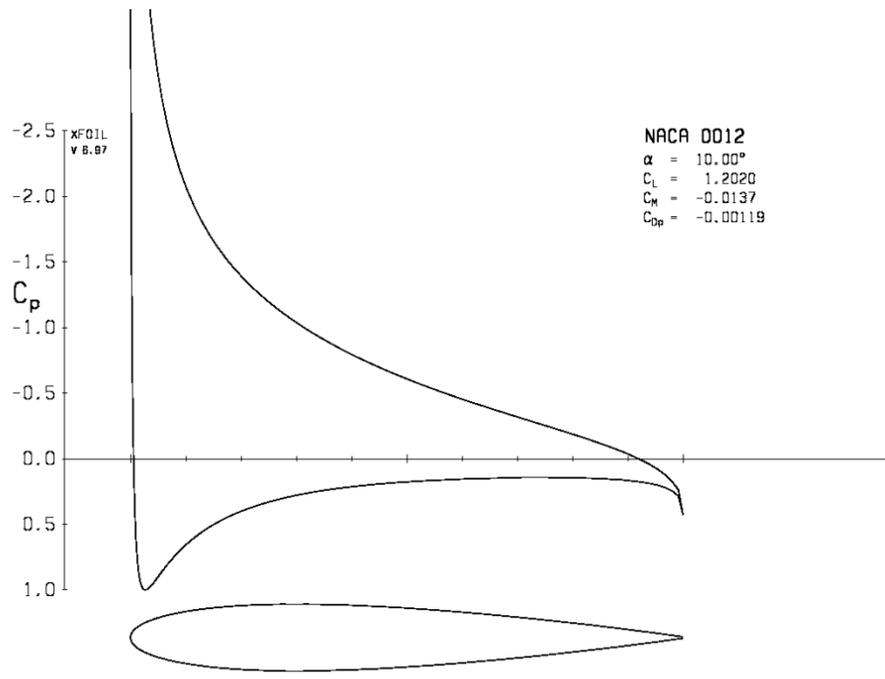
$$C_{m_{(x_r, y_r)}} = 0 \quad \Rightarrow \quad (x_r, y_r) = (x_{cp}, y_{cp})$$

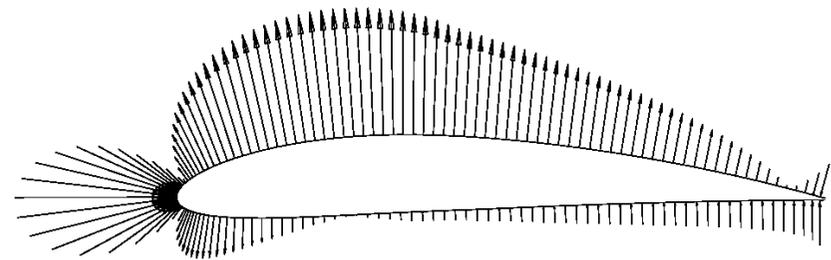
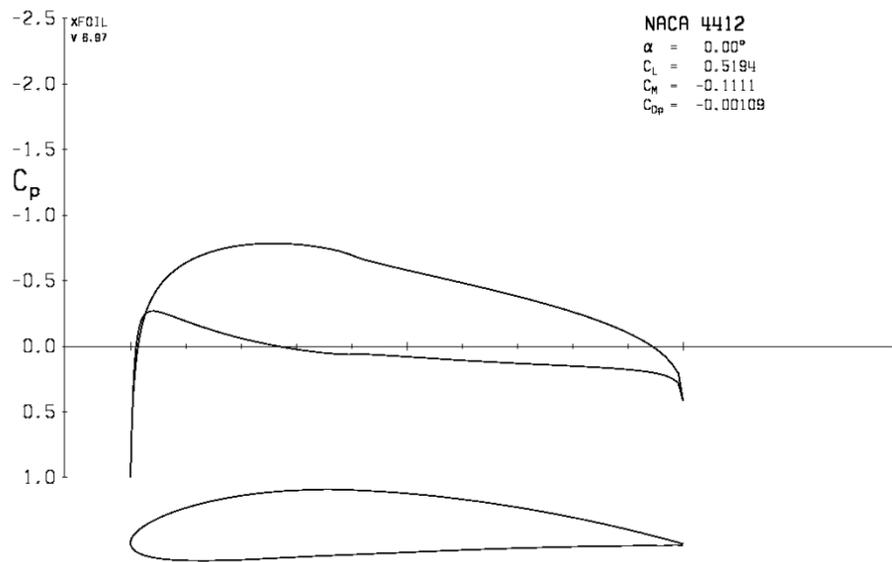
$$(x_{cp}/c) = (x_0/c) - \frac{C_{m_{(x_0, y_0)}}}{C_L}$$

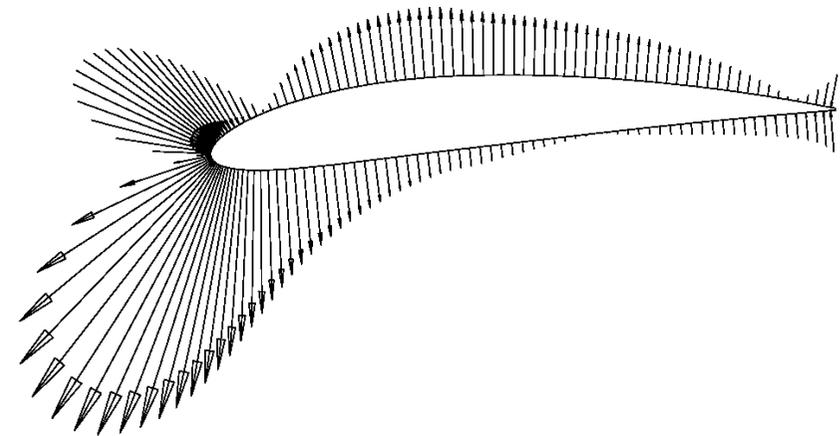
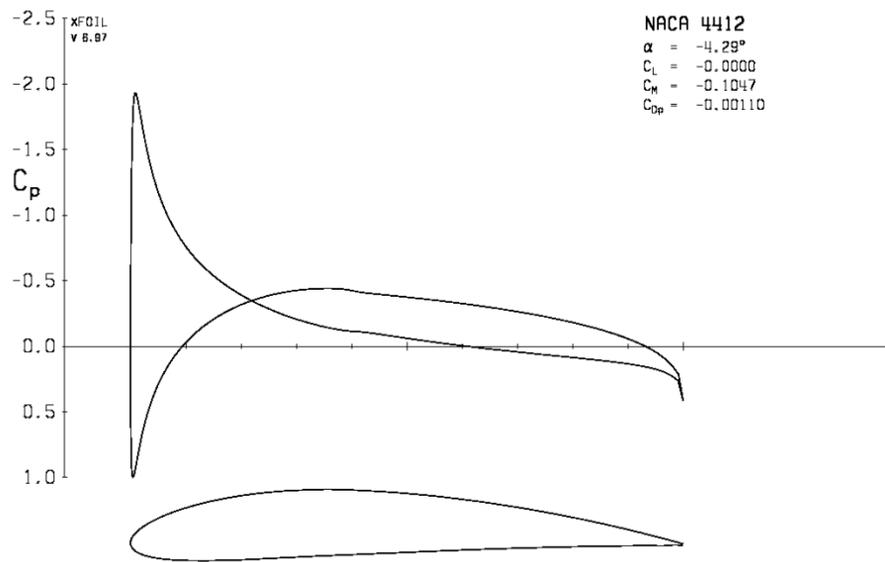


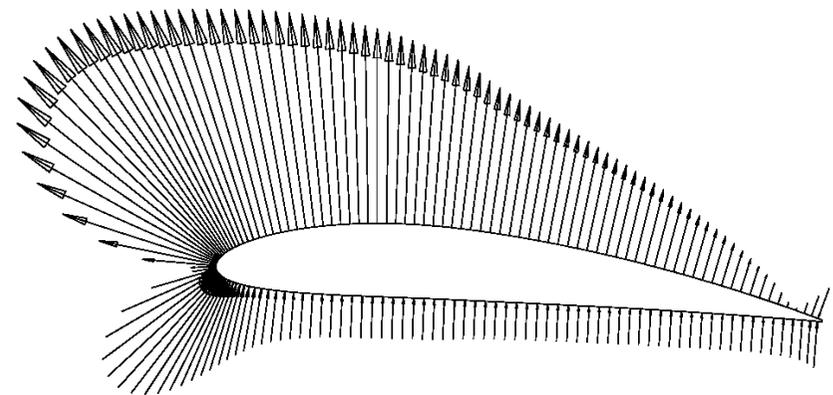
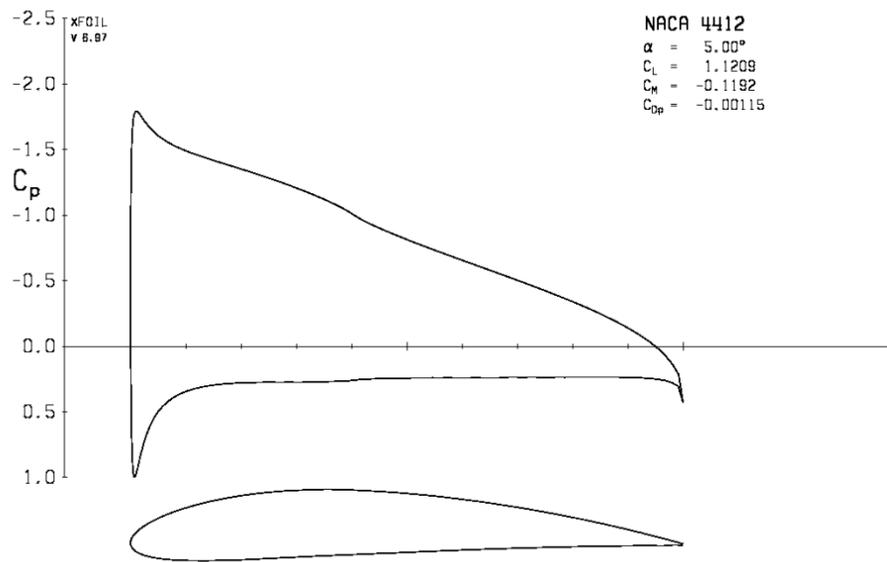


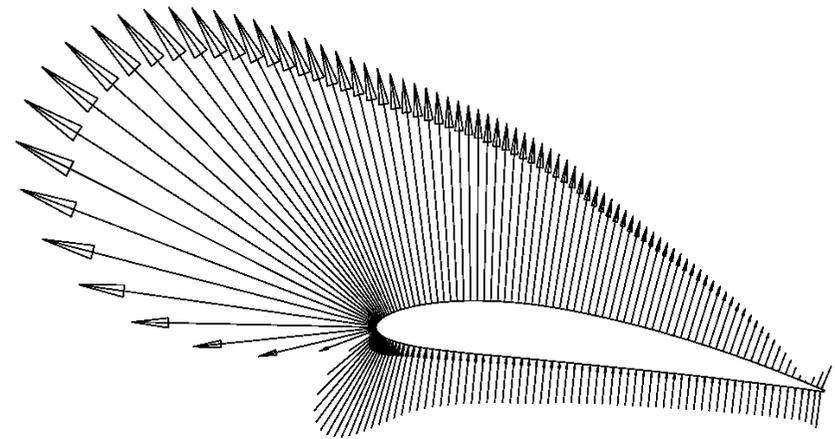
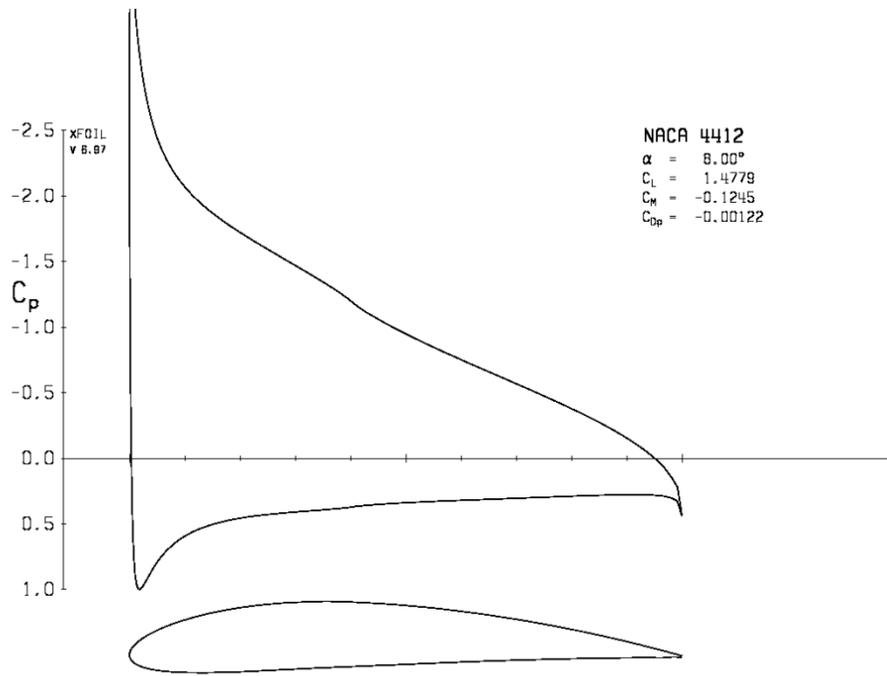












## BOUNDARY LAYER:

**x-momentum eq.:** 
$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = u_e \frac{du_e}{dx} + \frac{\partial}{\partial y} \left( \nu \frac{\partial u}{\partial y} \right) - \frac{\partial(\overline{u'v'})}{\partial y}$$

**y-momentum eq.:** 
$$\frac{\partial p}{\partial y} = 0, \quad p(y) = \text{const}$$

**continuity eq.:** 
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

Near the surface:  $u, v \ll U_\infty$  ( $\sim 0$ ), so:

Turbulent boundary layer, no pressure gradient:

$$\frac{\partial}{\partial y} \left( \nu \frac{\partial u}{\partial y} - \overline{u'v'} \right) = 0 \quad \text{lub:} \quad \nu \frac{\partial u}{\partial y} - \overline{u'v'} = \text{const} = \frac{\tau_w}{\rho}$$

def:

$$\frac{\tau_w}{\rho} = (V^*)^2, \quad \frac{u}{V^*} = u^+, \quad \frac{y}{\nu/V^*} = y^+$$

near the surface:

$$\overline{u'v'} \ll \nu \frac{\partial u}{\partial y}$$

so:

$$\nu \frac{\partial u}{\partial y} = \frac{\tau_w}{\rho} = (V^*)^2$$

$$u = \frac{(V^*)^2}{\nu} \cdot y \quad \rightarrow \quad u^+ = y^+$$

***LAMINAR (VISCIOUS) SUBLAYER***

above laminar sublayer:

$$\overline{u'v'} \gg \nu \frac{\partial u}{\partial y}$$

so:

$$\overline{u'v'} = (V^*)^2$$

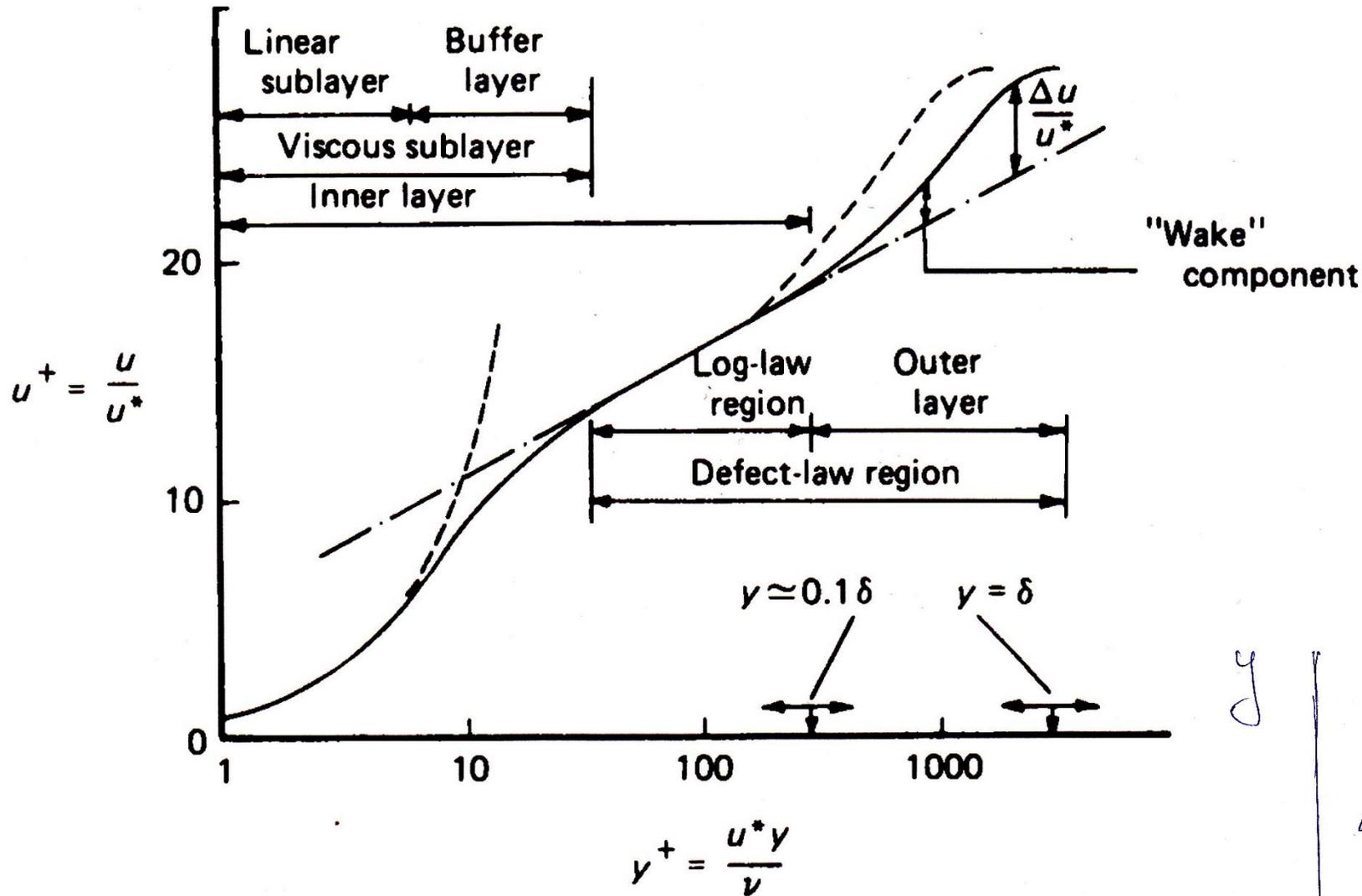
turbulence model (mixing length):

$$\nu_T = l_m^2 \left| \frac{\partial u}{\partial y} \right|, \quad l_m = \kappa \cdot y$$

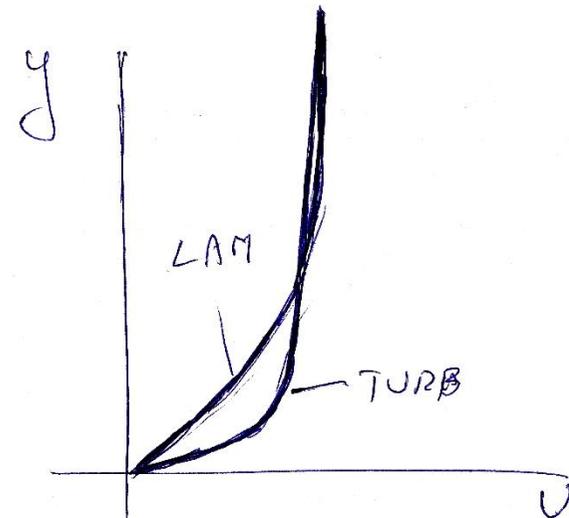
$$l_m^2 \left( \frac{\partial u}{\partial y} \right)^2 = \frac{\tau_w}{\rho} = (V^*)^2, \quad l_m \left( \frac{\partial u}{\partial y} \right) = V^*, \quad \frac{\partial u}{\partial y} = \frac{V^*}{\kappa \cdot y} \rightarrow u = \frac{V^*}{\kappa} \ln y + C_1$$

$$u^+ = \frac{\ln y^+}{\kappa} + C \quad \kappa = 0.41, \quad C = 6.1$$

J.J.Bertin, M.L.Smith – Aerodynamics for Engineers, Printice-Hall International, Inc.,  
1998



INTEGRAL RELATIONS FOR BOUNDARY LAYER (ORI



$$\frac{\partial(uu)}{\partial x} + \frac{\partial(uv)}{\partial y} = u_e \frac{du_e}{dx} + v \frac{\partial^2 u}{\partial y^2} \quad (1)$$

$$\frac{\partial(u_e u)}{\partial x} + \frac{\partial(u_e v)}{\partial y} = u \frac{du_e}{dx} \quad (2)$$

(2)-(1)

$$\frac{\partial((u_e - u)u)}{\partial x} + \frac{\partial((u_e - u)v)}{\partial y} = (u_e - u) \frac{du_e}{dx} - v \frac{\partial^2 u}{\partial y^2}$$

$$\int_0^{\delta} \frac{\partial((u_e - u)u)}{\partial x} dy + \int_0^{\delta} \frac{\partial((u_e - u)v)}{\partial y} dy + \int_0^{\delta} (u_e - u) \frac{du_e}{dx} dy = -v \int_0^{\delta} \frac{\partial^2 u}{\partial y^2} dy$$

$$\frac{d}{dx} \int_0^{\delta} (u_e - u)u dy - [(u_e - u)u]_{y=\delta} \frac{d\delta}{dx} +$$

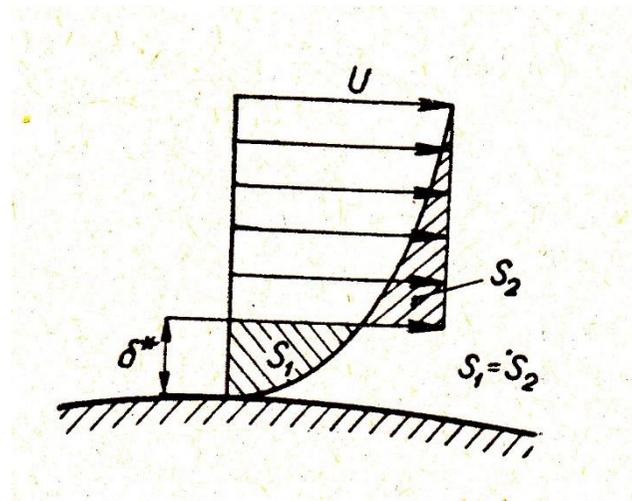
$$((u_e - u)v)|_0^{\delta} + \frac{du_e}{dx} \int_0^{\delta} (u_e - u) dy = -v \frac{\partial u}{\partial y} \Big|_0^{\delta}$$

$$\frac{d}{dx} u_e^2 \int_0^{\delta} \frac{u}{u_e} \left(1 - \frac{u}{u_e}\right) dy + \frac{du_e}{dx} u_e \int_0^{\delta} \left(1 - \frac{u}{u_e}\right) dy - \frac{\tau_w}{\rho} = 0$$

$$\delta^* = \int_0^{\delta} \left(1 - \frac{u}{u_e}\right) dy$$

$$\theta = \int_0^{\delta} \frac{u}{u_e} \left(1 - \frac{u}{u_e}\right) dy$$

$$C_f = \frac{\tau_w}{\rho_e u_e^2 / 2}$$



$$H = \frac{\delta^*}{\theta}$$

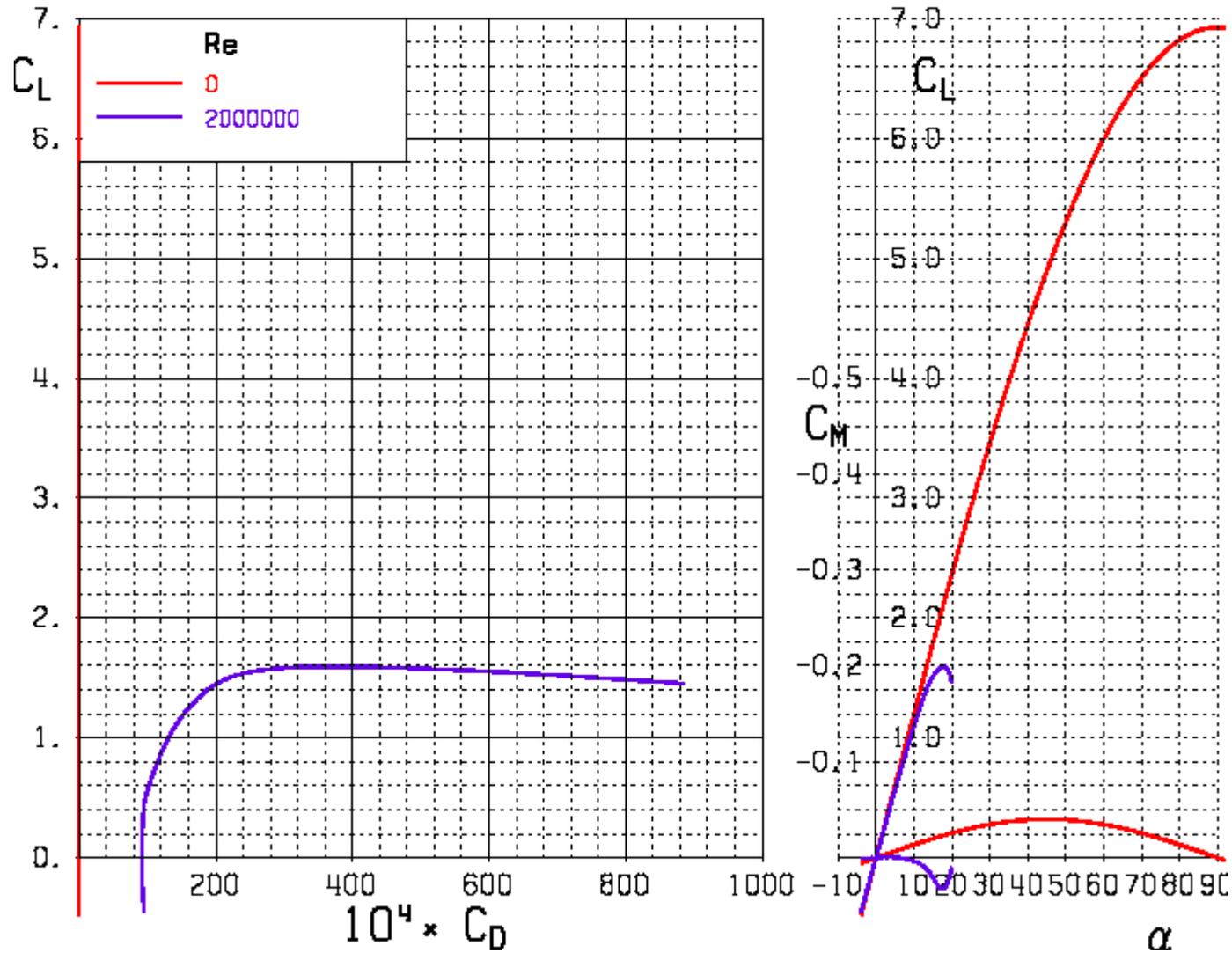
$$\frac{d}{dx}(u_e^2 \theta) + \frac{du_e}{dx} u_e \delta^* - \frac{\tau_w}{\rho} = 0$$

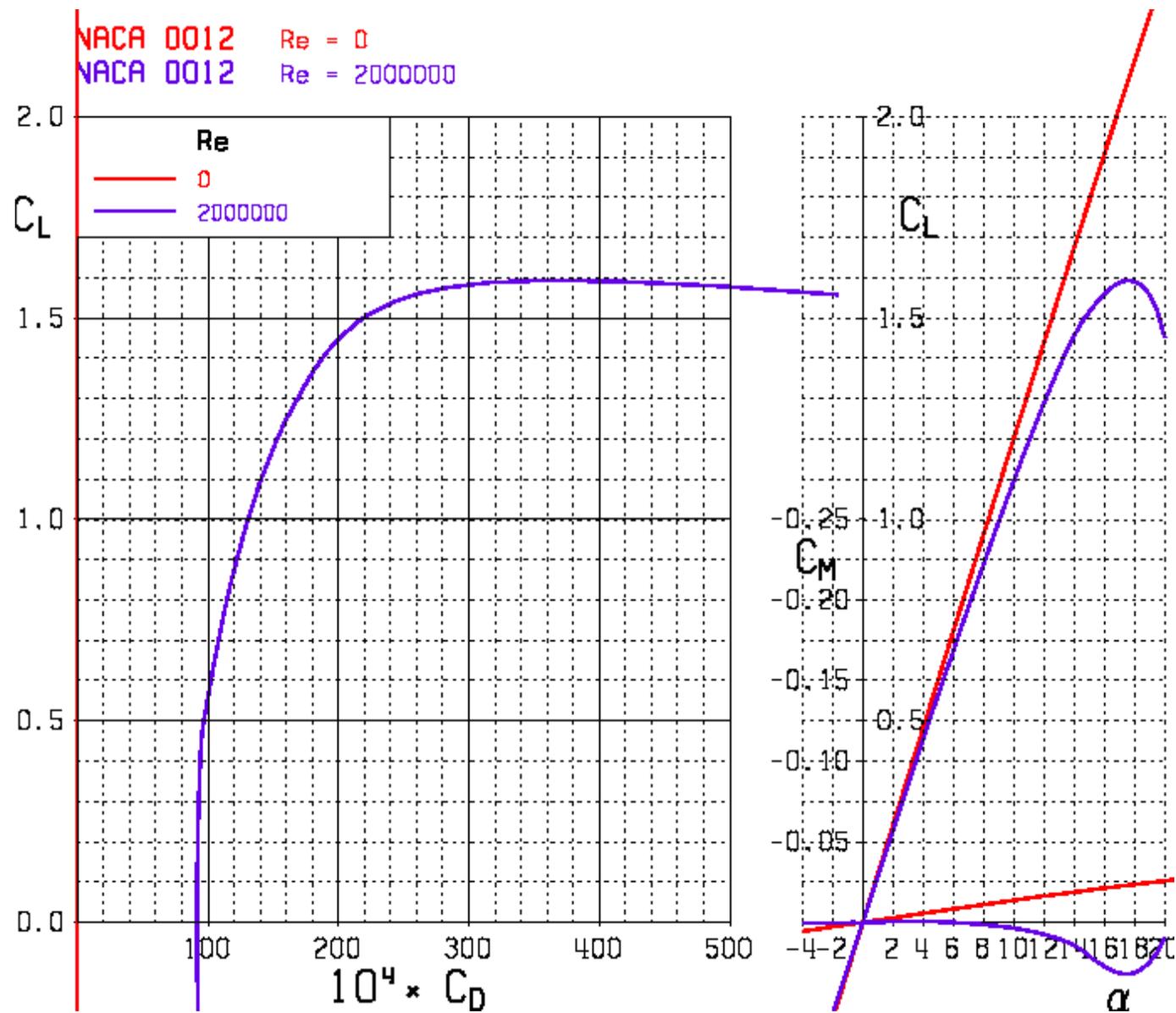
$$u_e^2 \frac{d\theta}{dx} + 2u_e \frac{du_e}{dx} \theta + u_e \frac{du_e}{dx} \delta^* - u_e^2 \frac{\tau_w}{\rho_e u_e^2} = 0$$

$$\frac{d\theta}{dx} + \frac{\theta}{u_e} \frac{du_e}{dx} \left(2 + \frac{\delta^*}{\theta}\right) - \frac{1}{2} C_f = 0$$

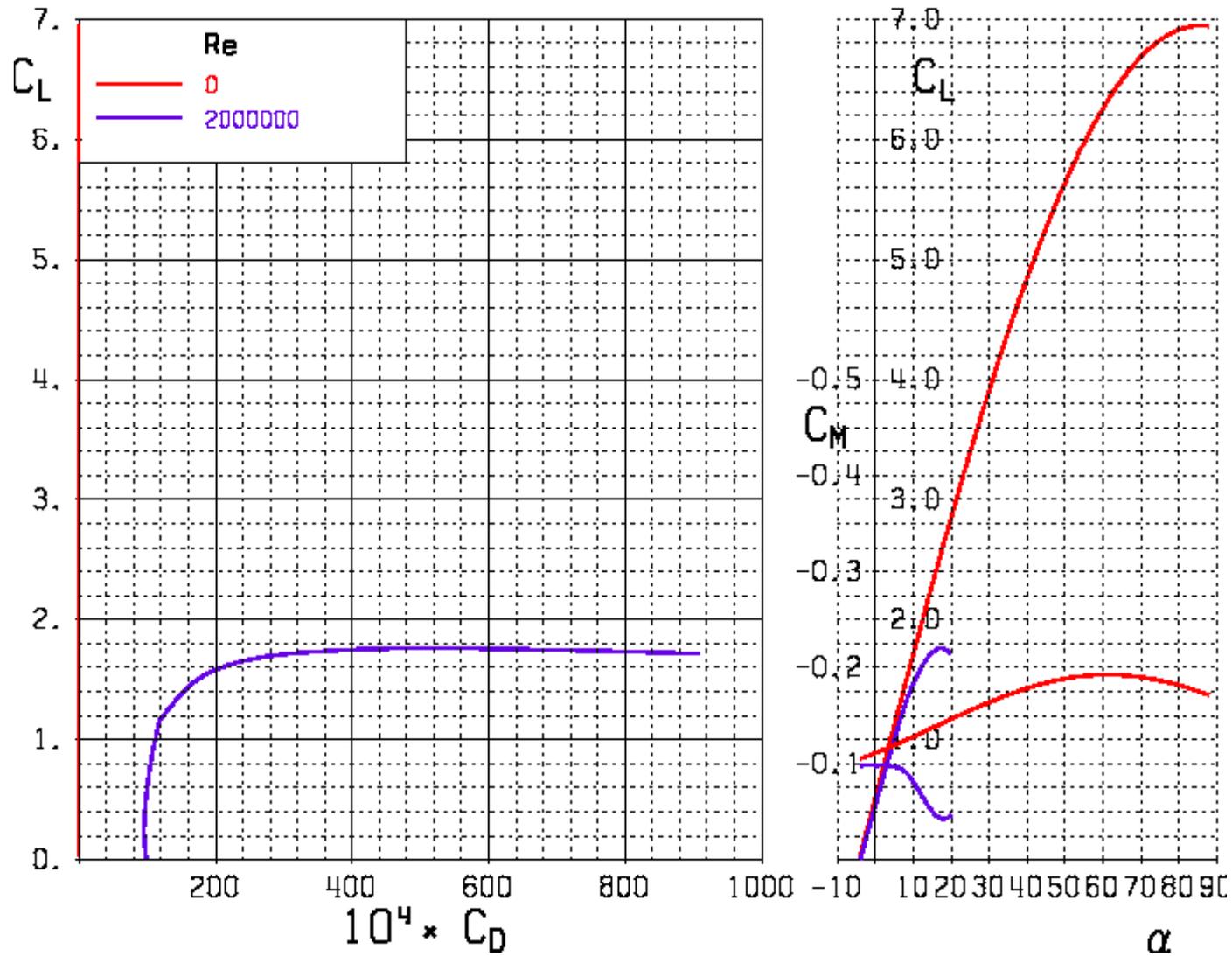
$$\frac{d\theta}{dx} + \frac{\theta}{u_e} \frac{du_e}{dx} (2 + H) - \frac{1}{2} C_f = 0$$

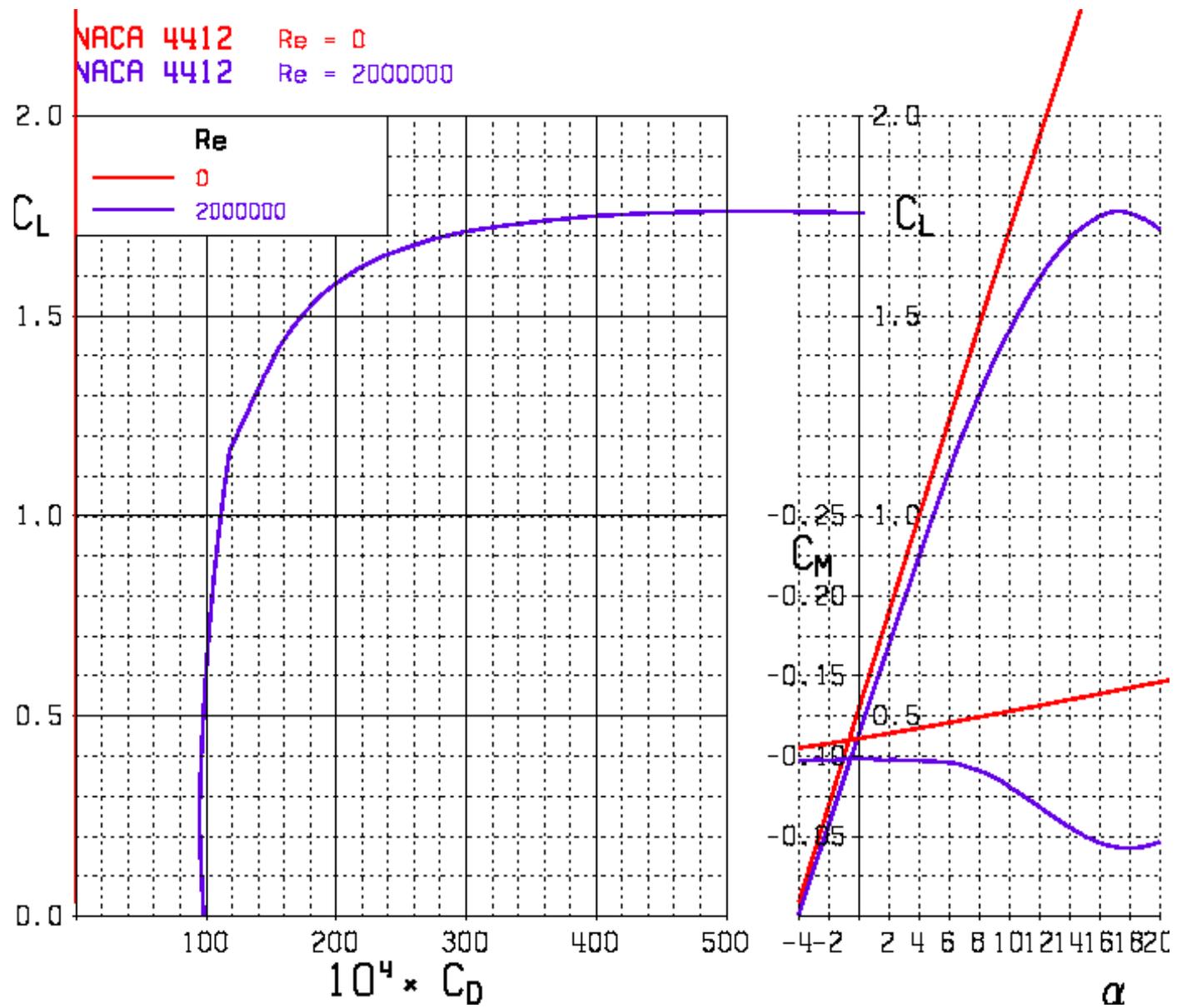
NACA 0012  $Re = 0$   
 NACA 0012  $Re = 2000000$

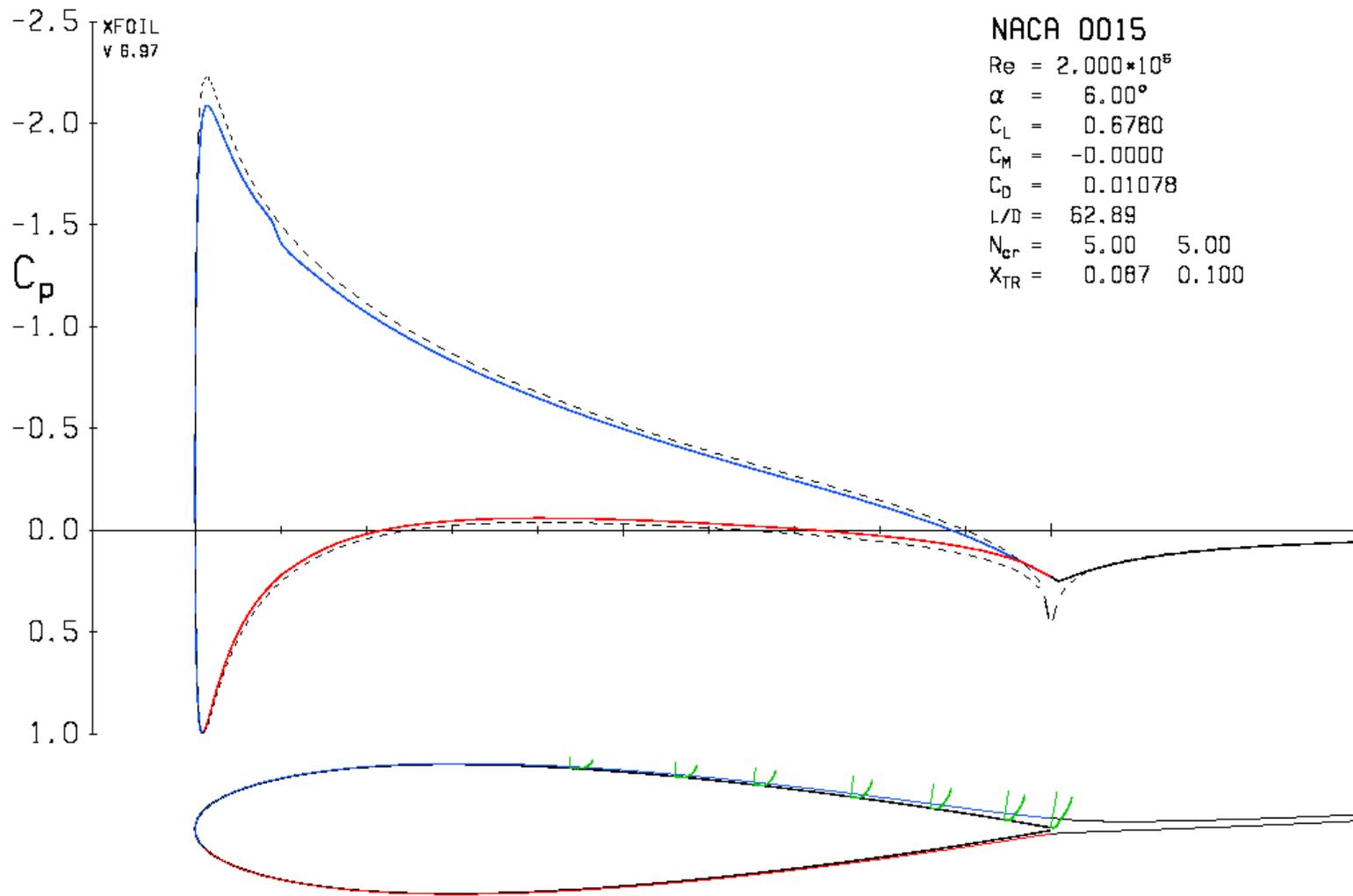


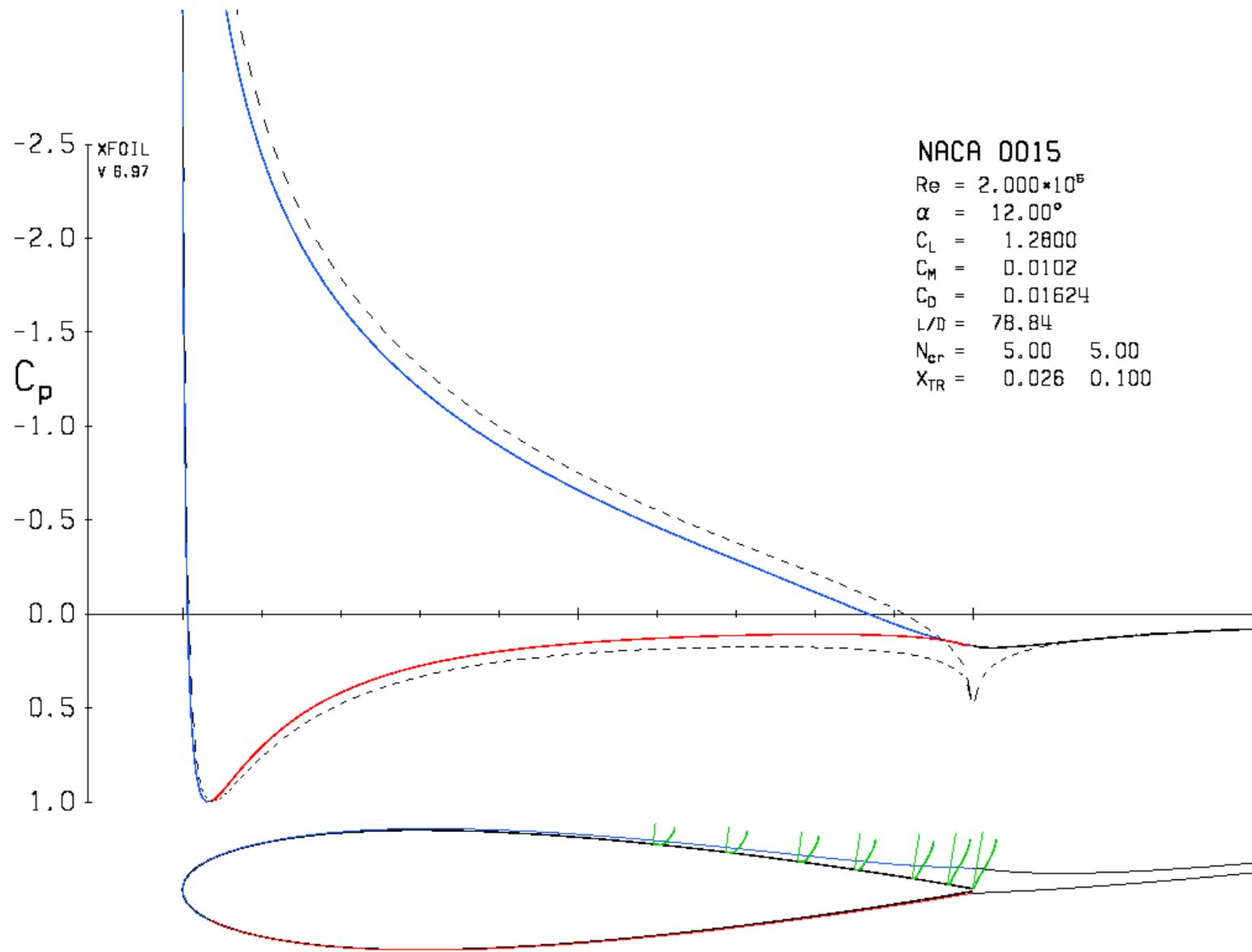


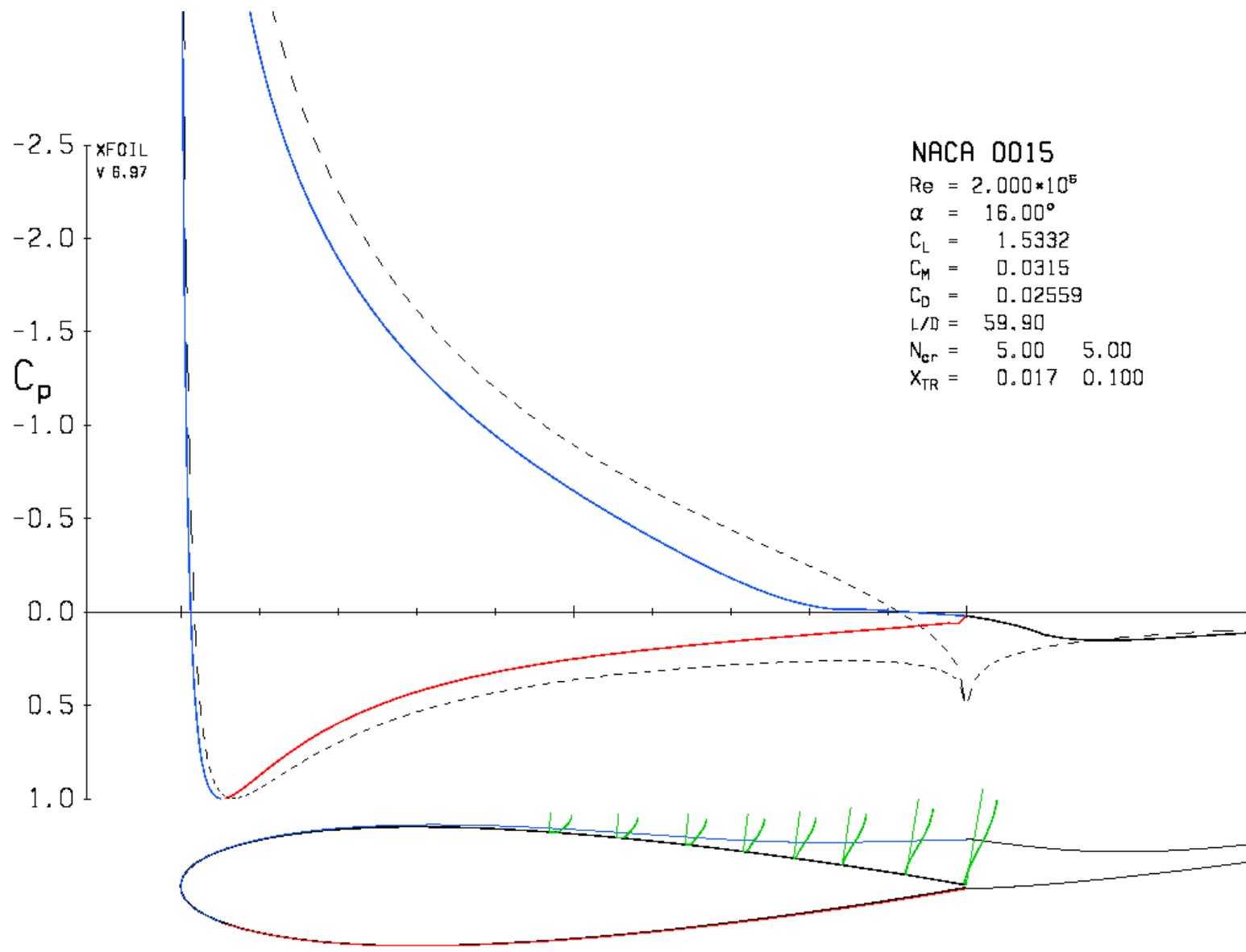
NACA 4412  $Re = 0$   
 NACA 4412  $Re = 2000000$

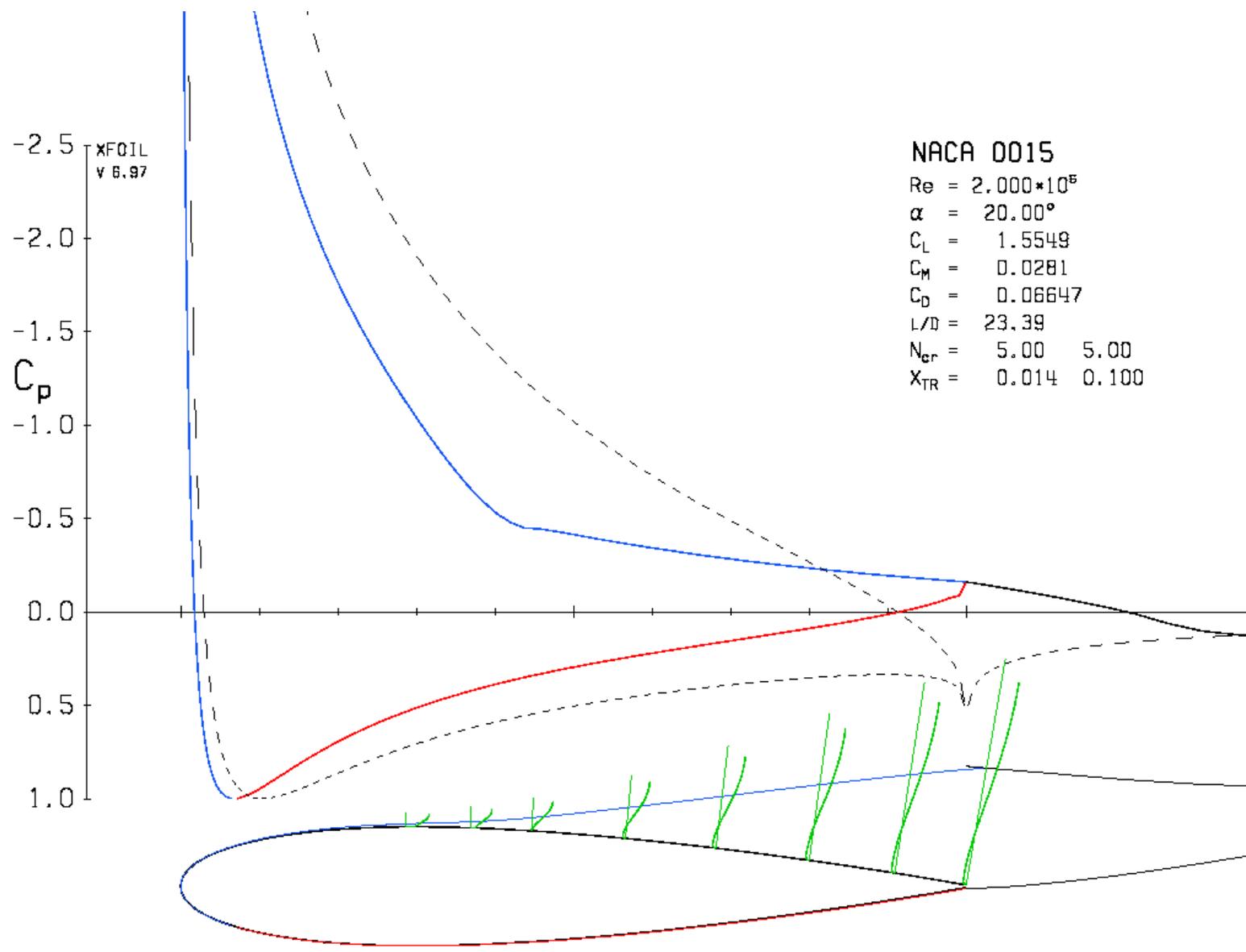


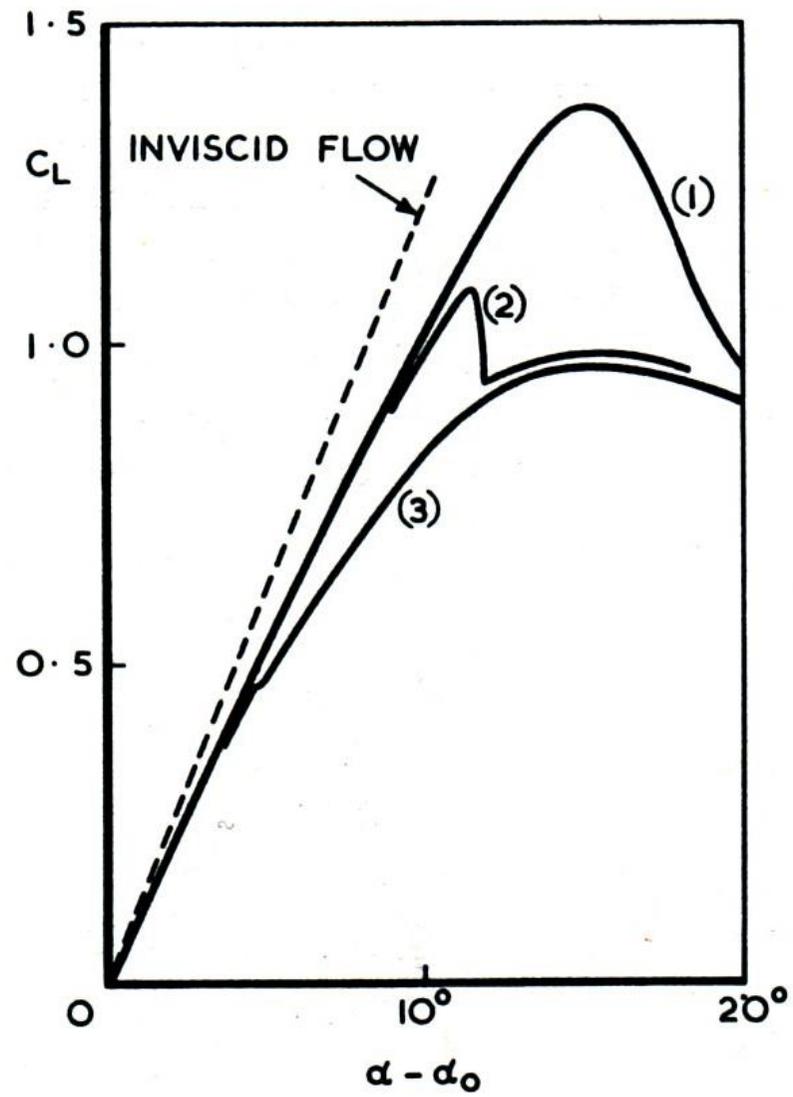




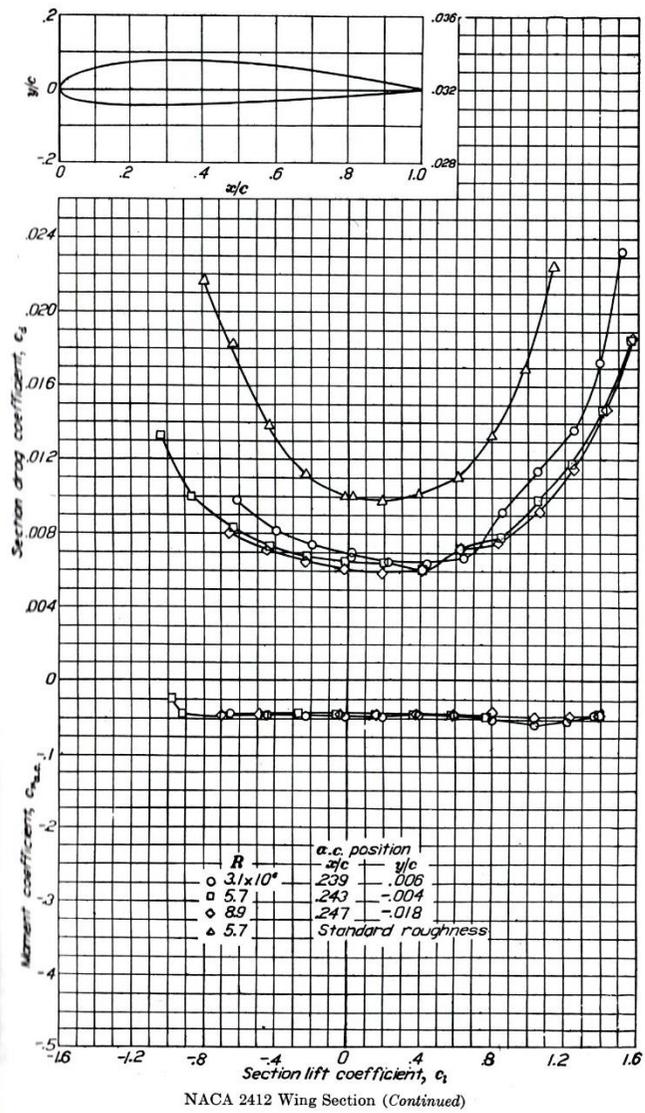
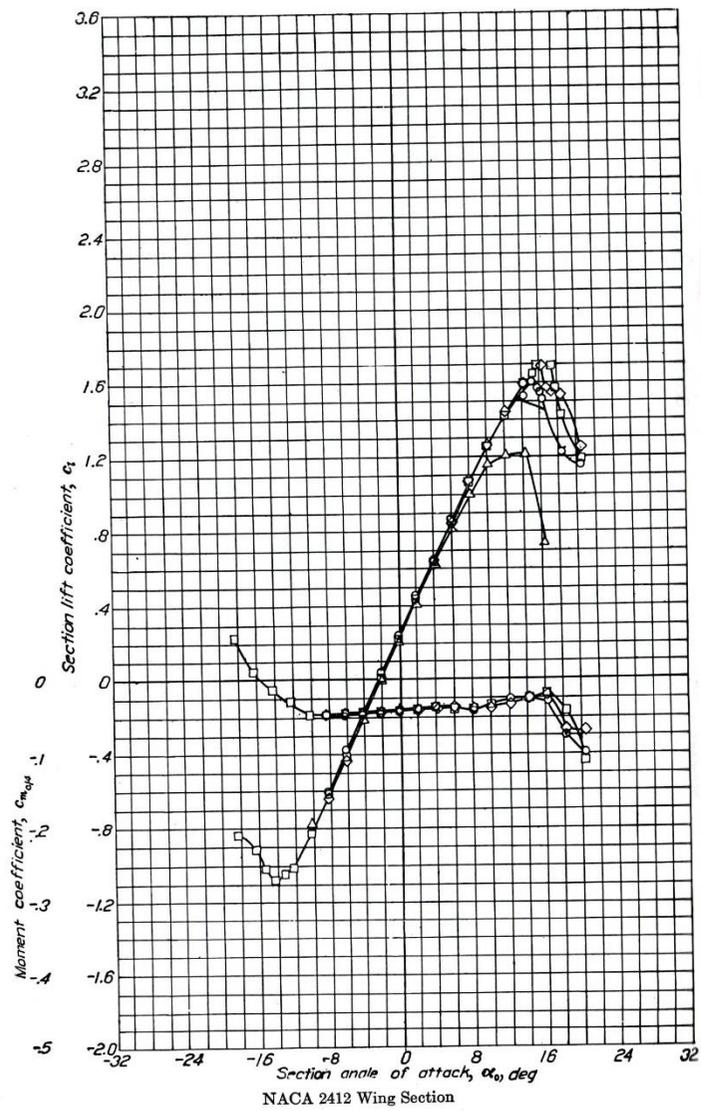


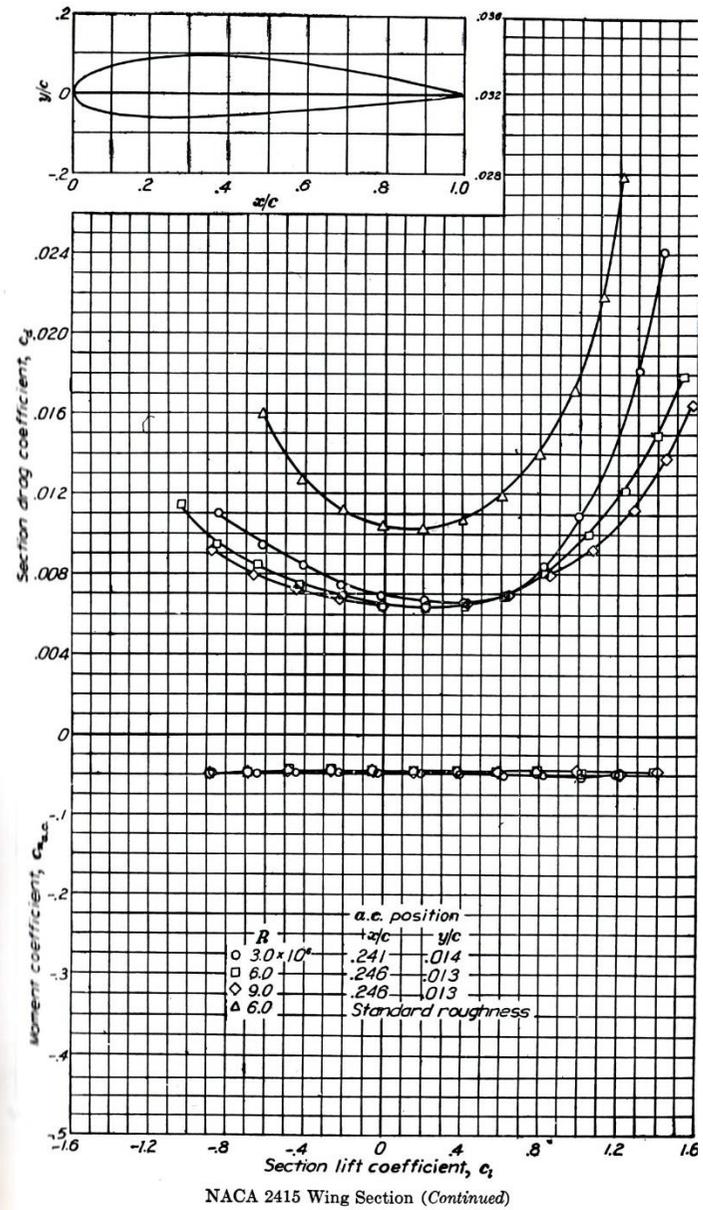
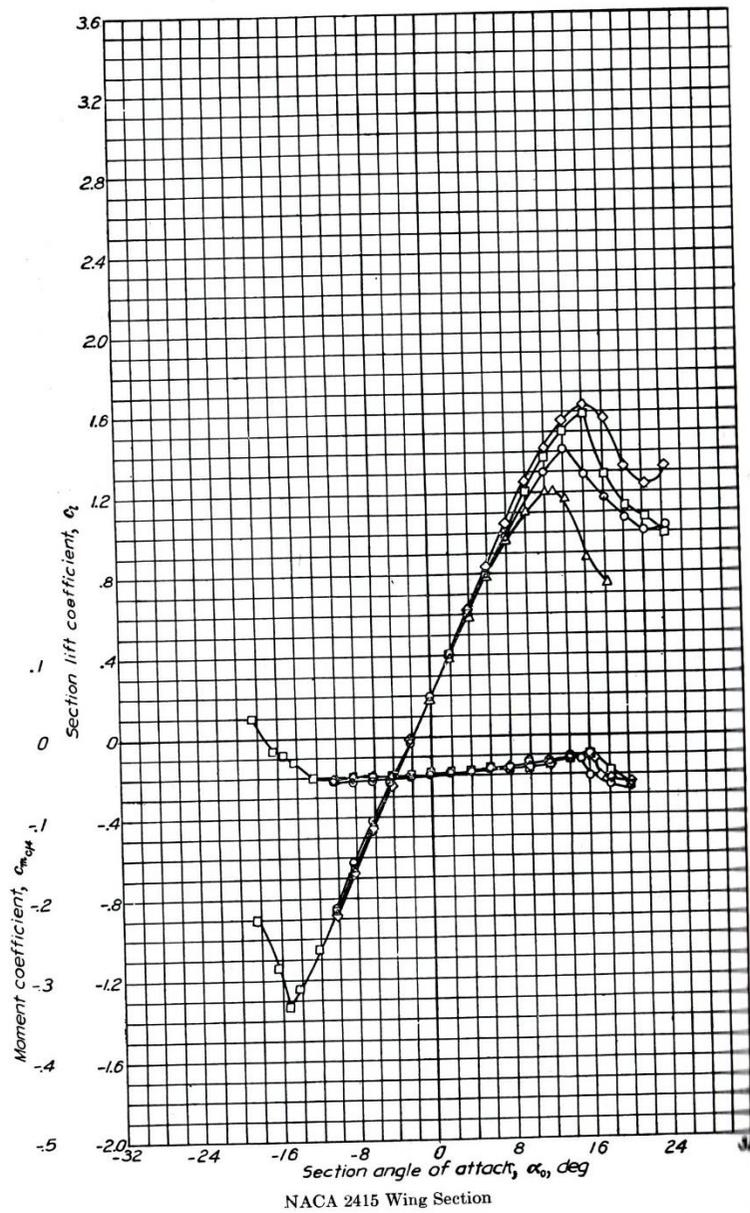


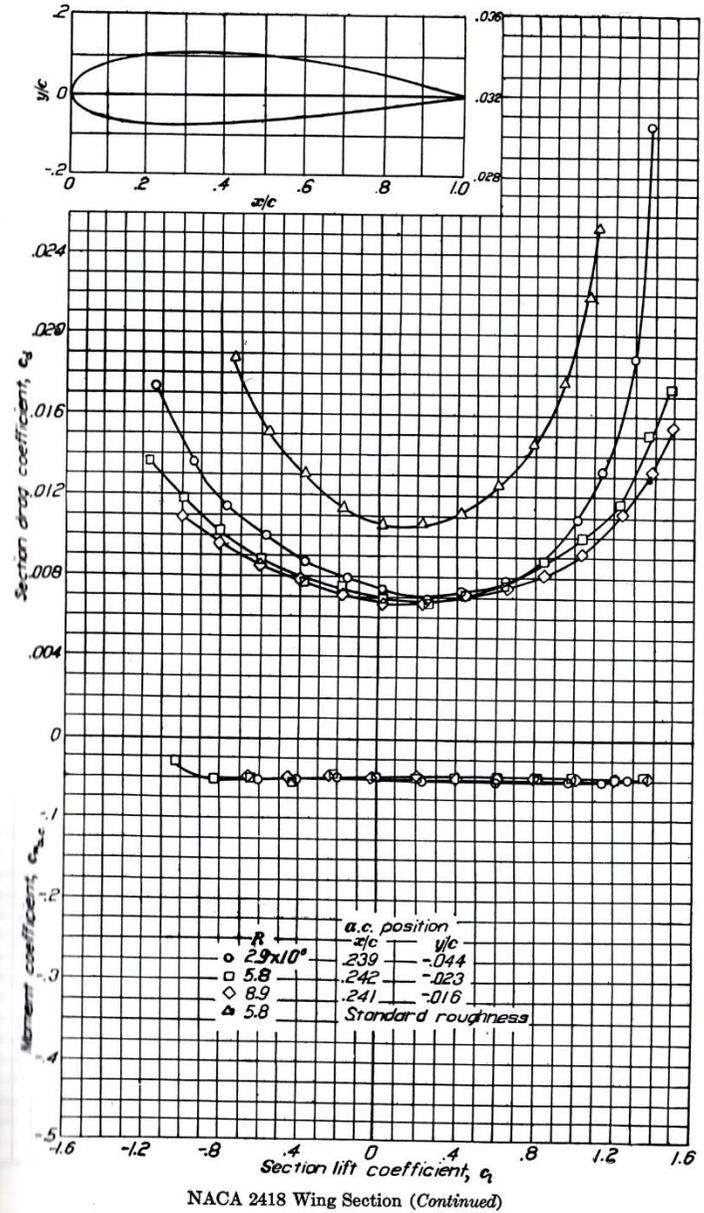
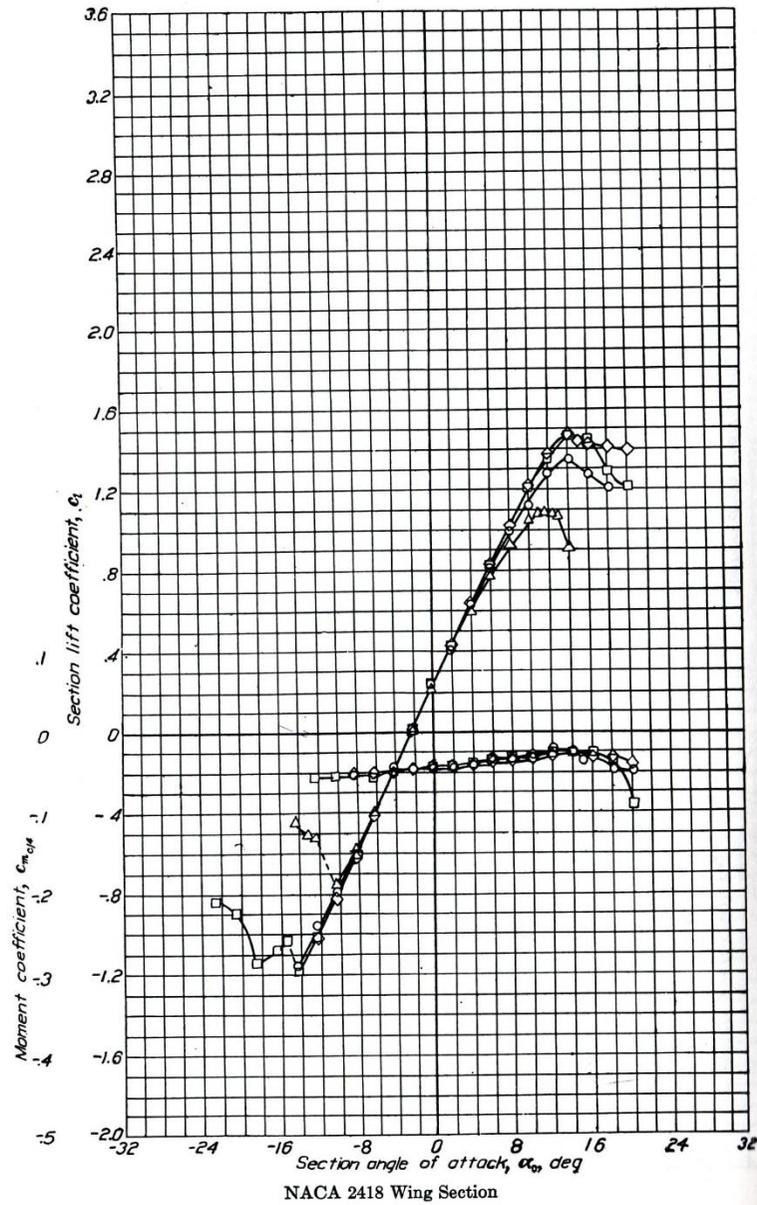


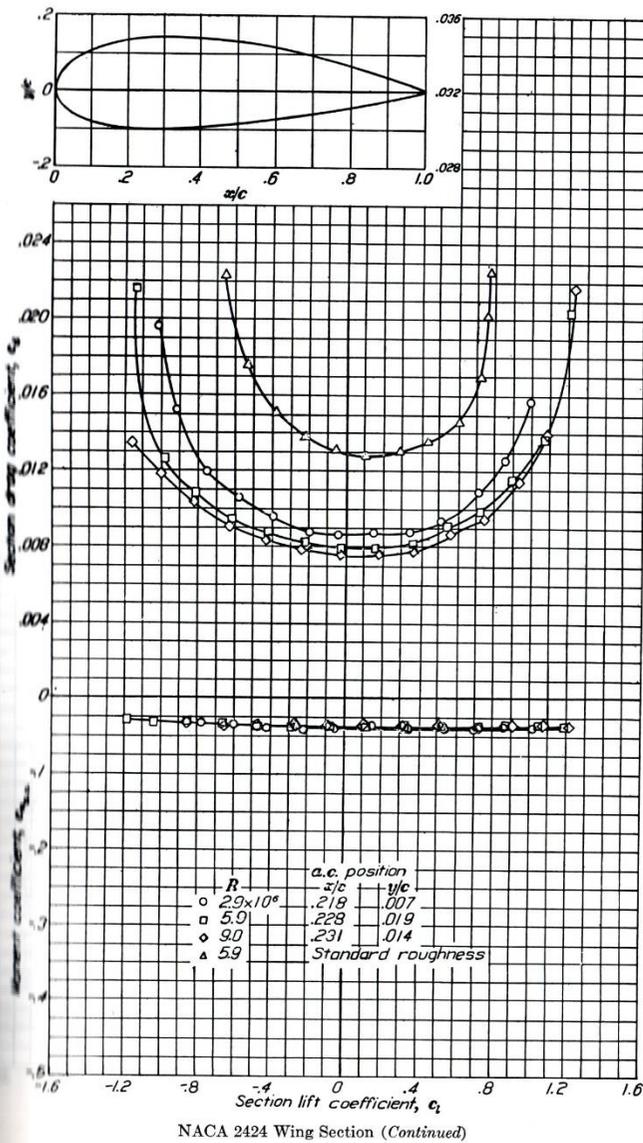
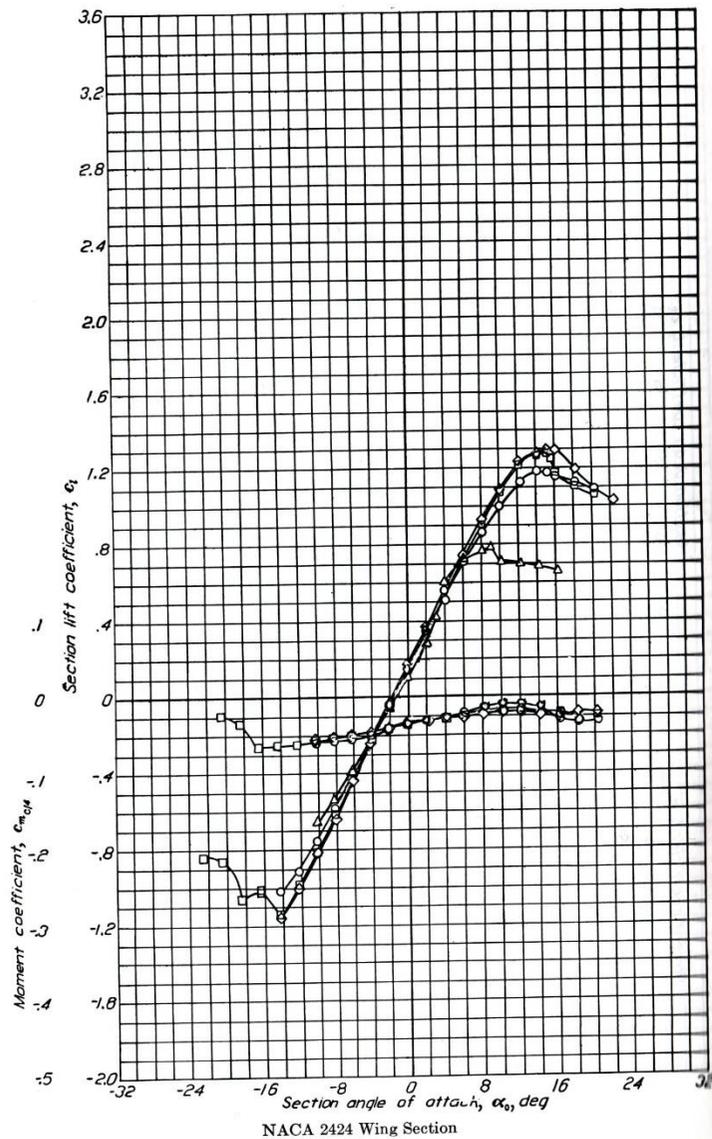


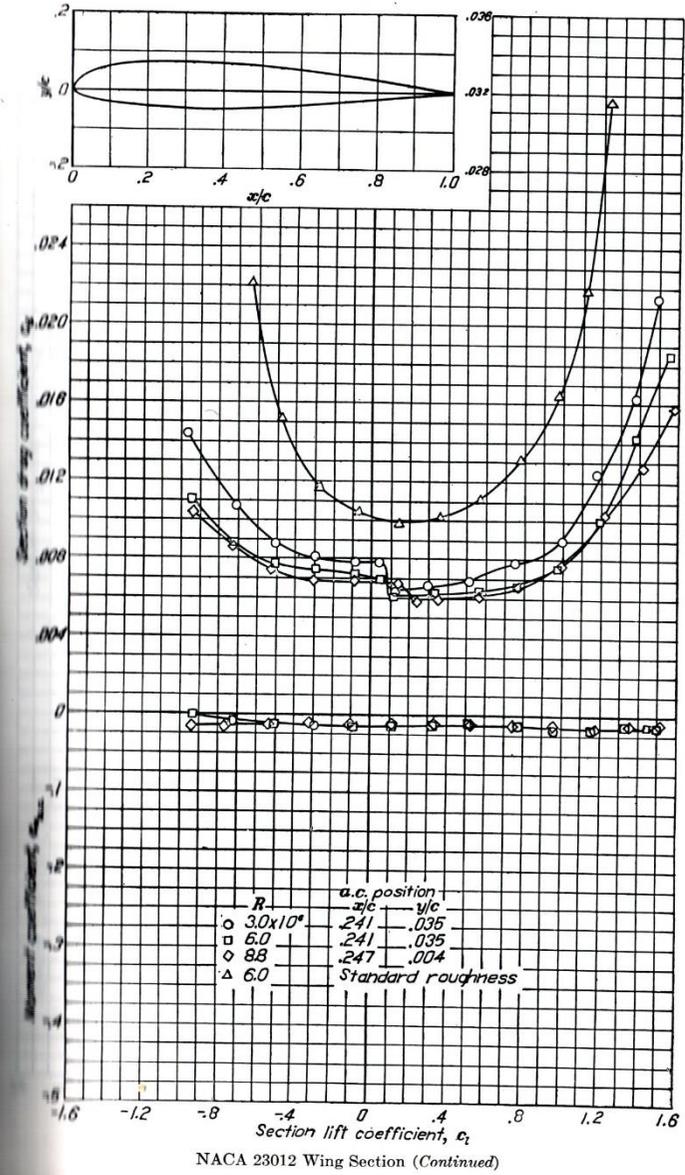
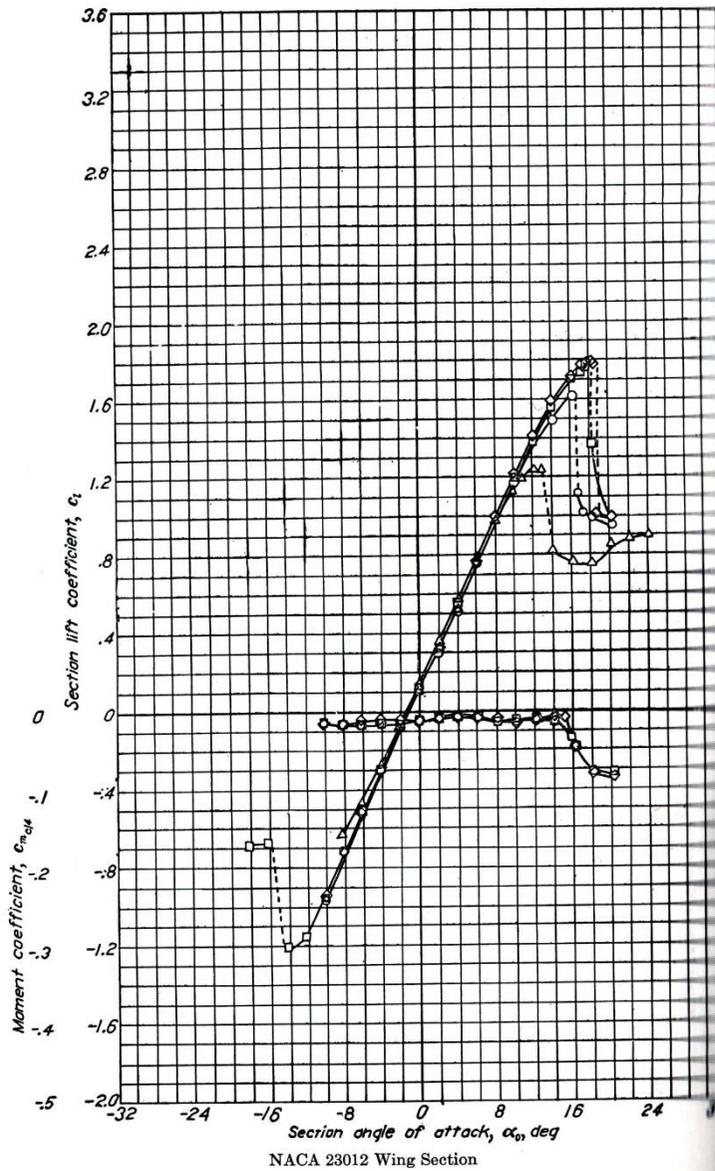
I.H.Abbott, A.E.von Doenhoff, Theory of Wing Sections, Dover Publications Inc., 1958 (presentation pp. 37-53)

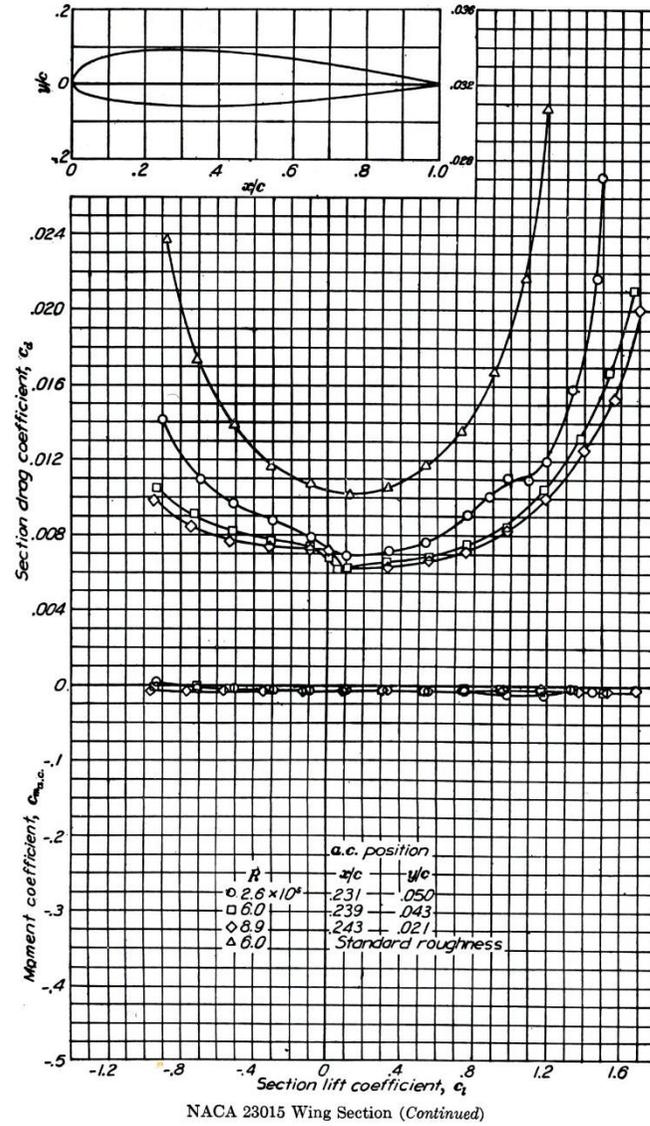
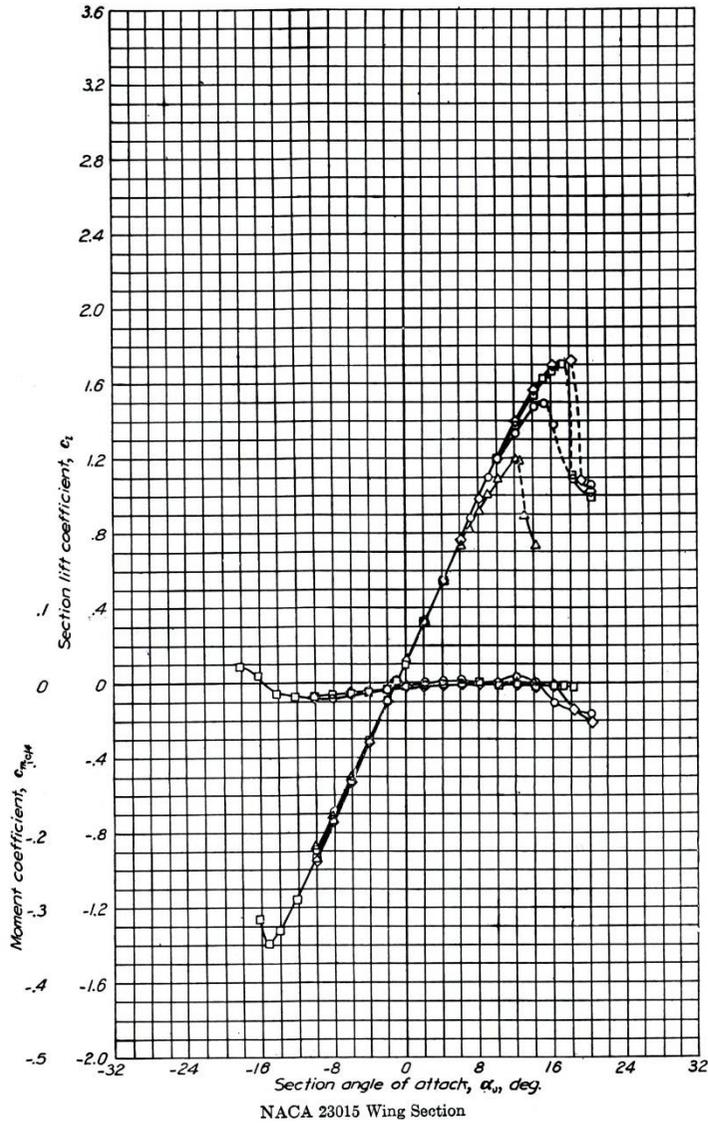


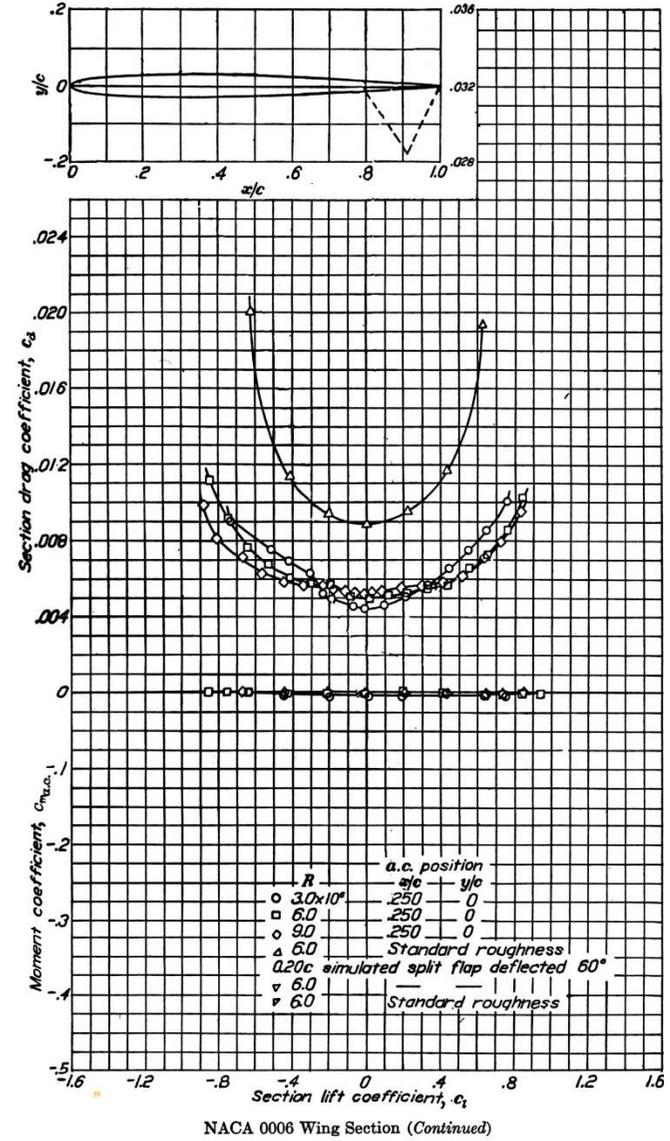
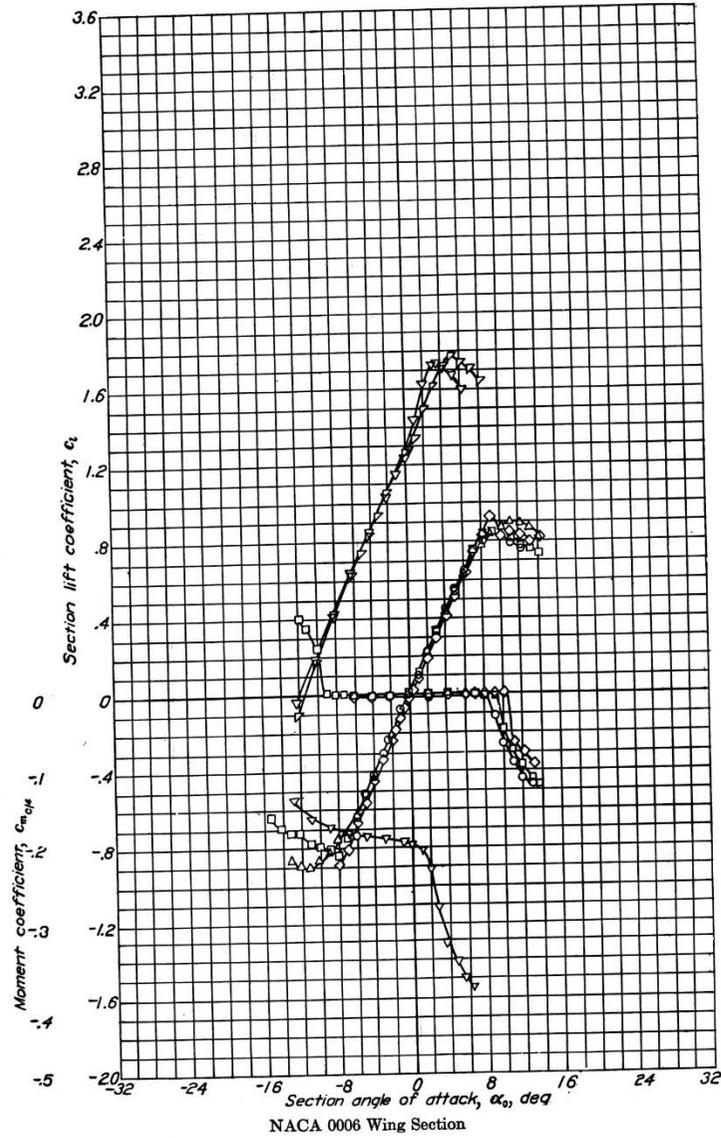












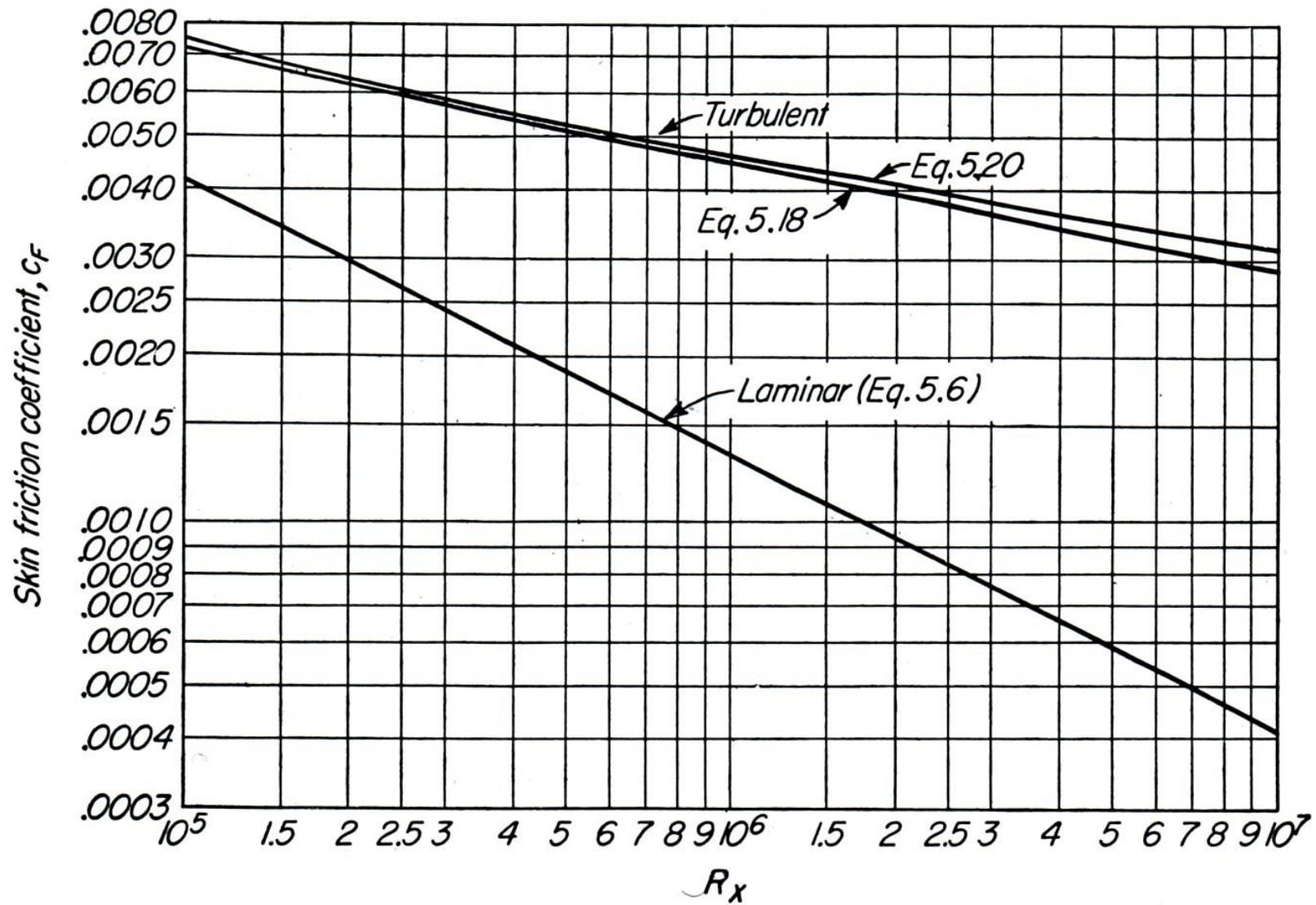
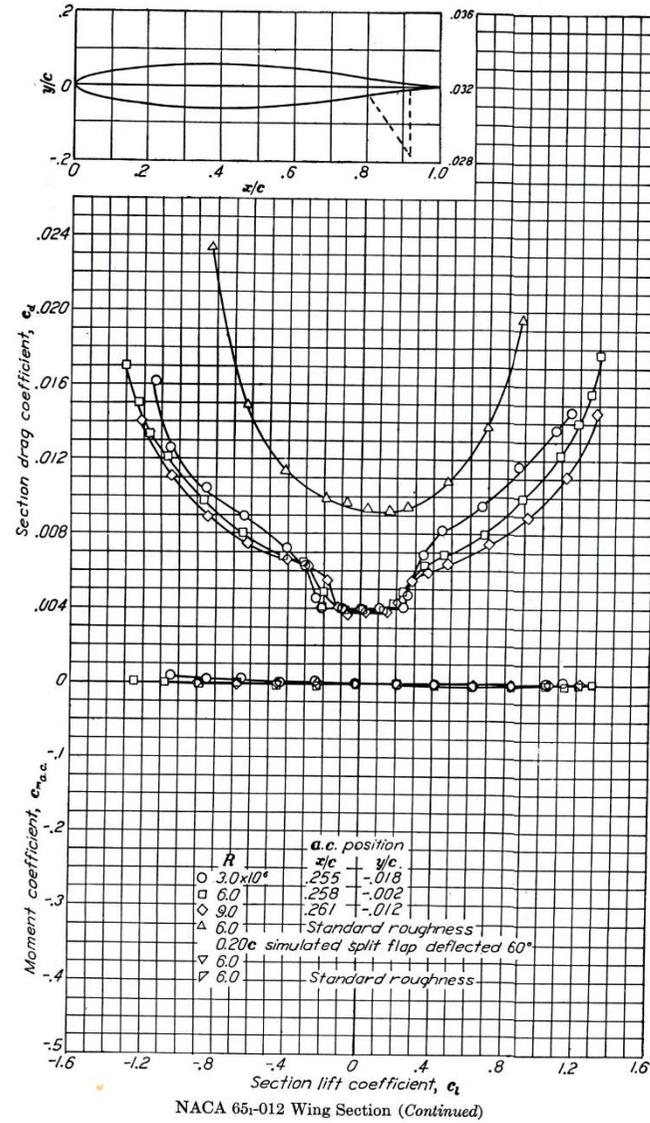
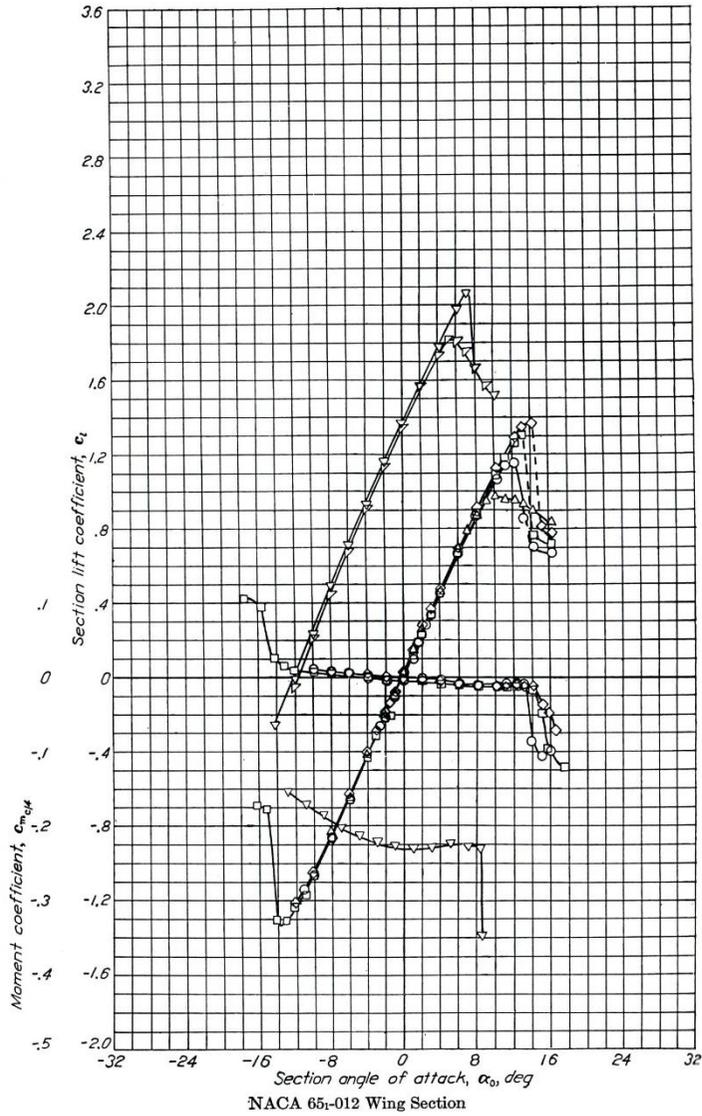
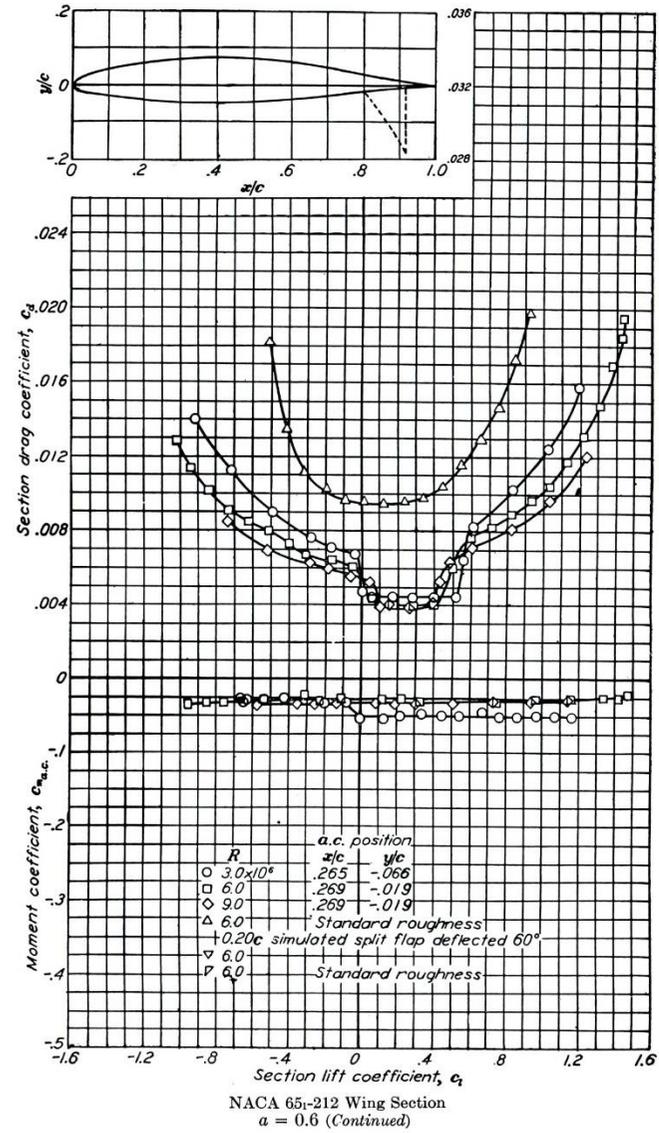
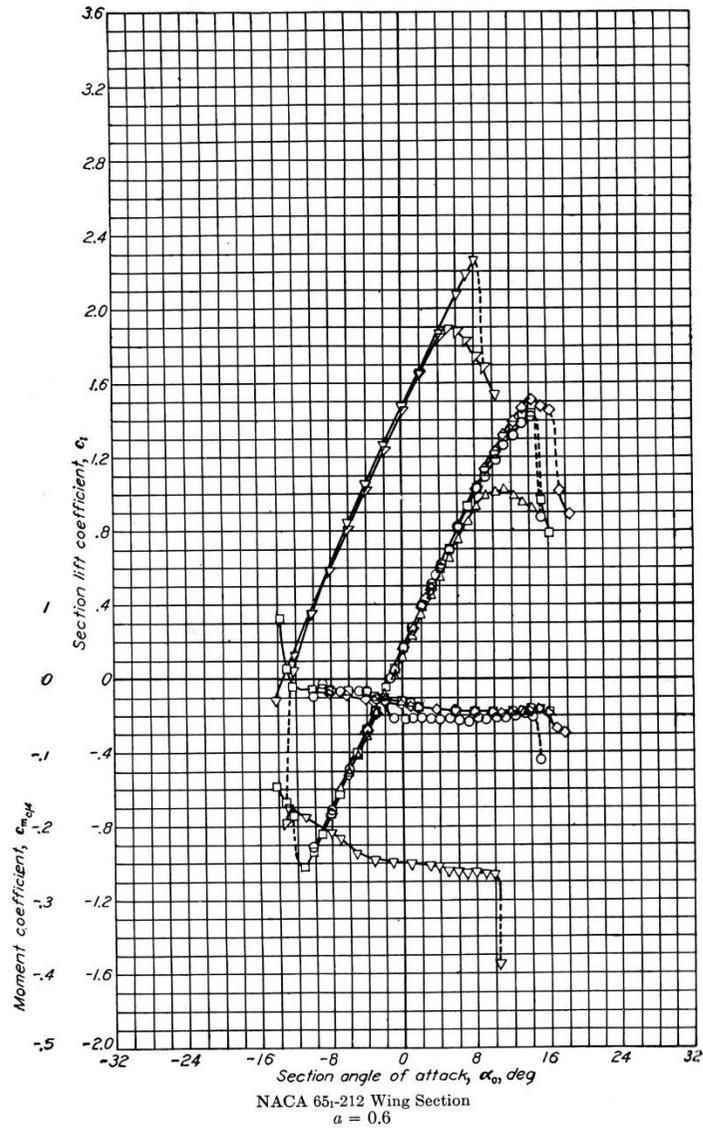
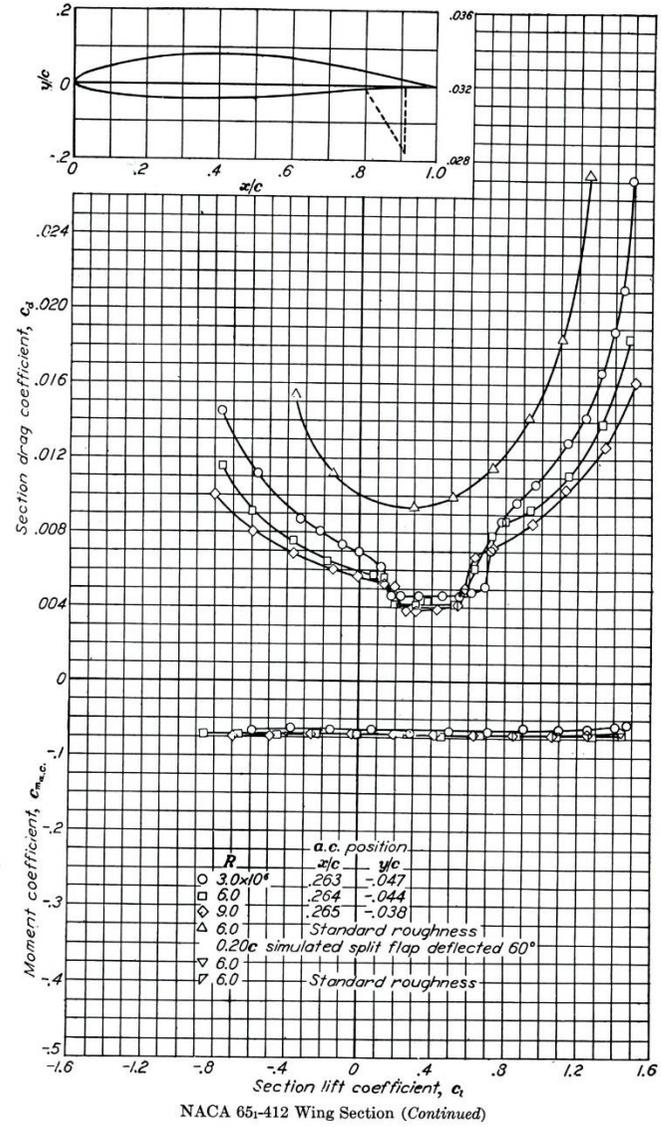
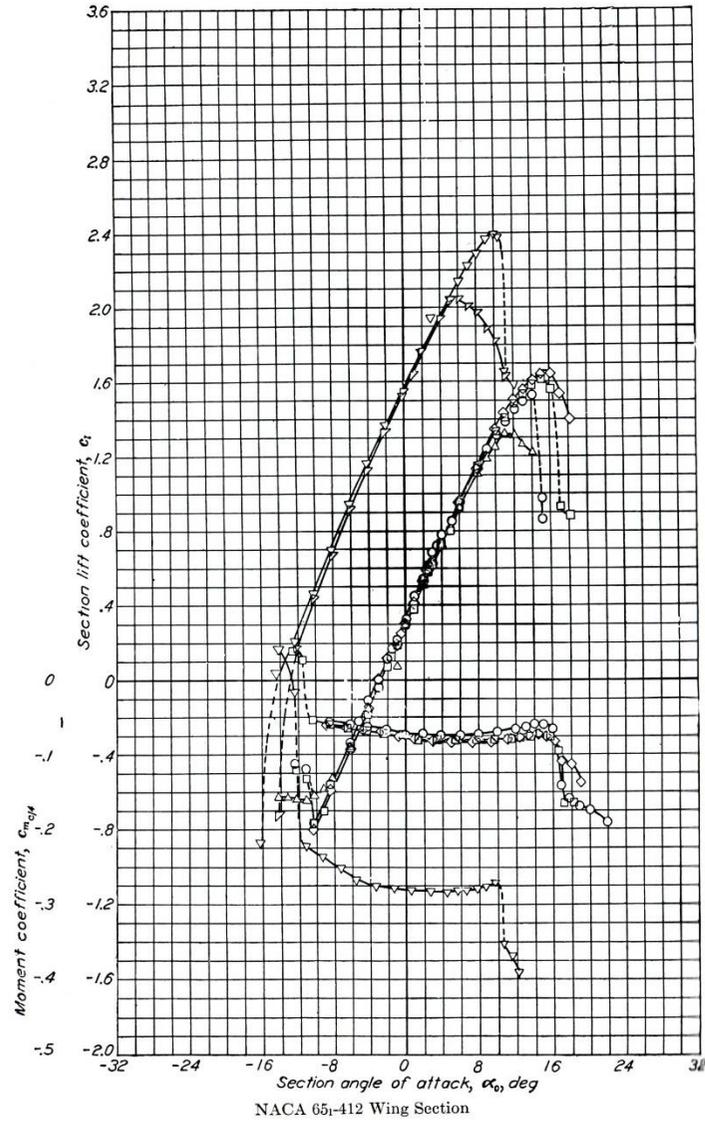
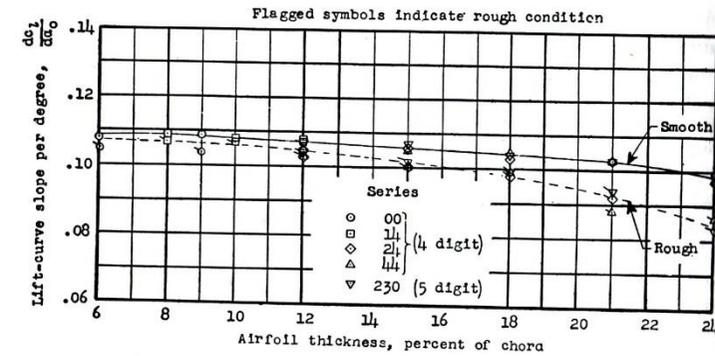


FIG. 46. Laminar and turbulent skin-friction coefficients for one side of a flat plate.

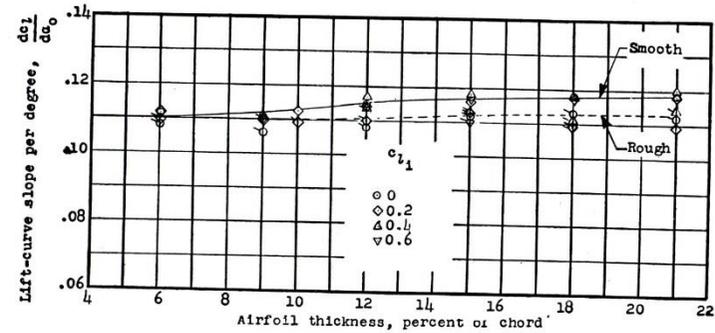




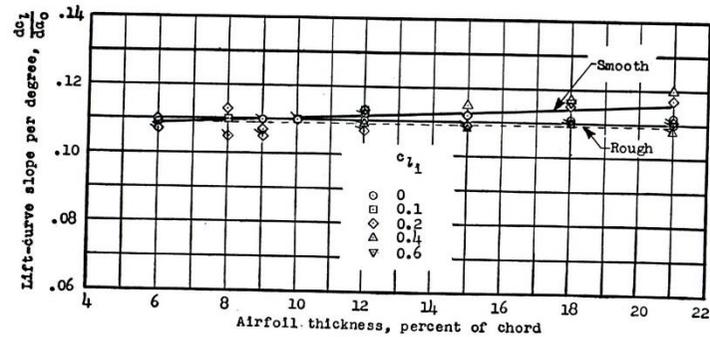




(a) NACA four- and five-digit series.

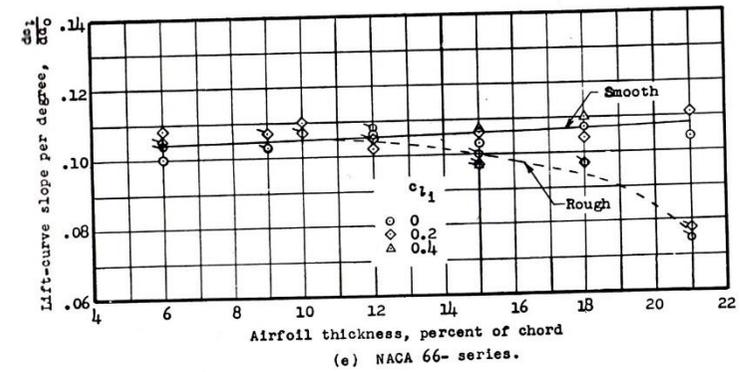
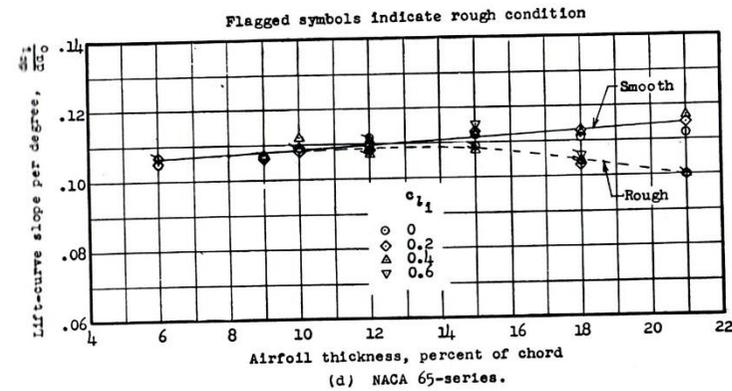


(b) NACA 63- series.

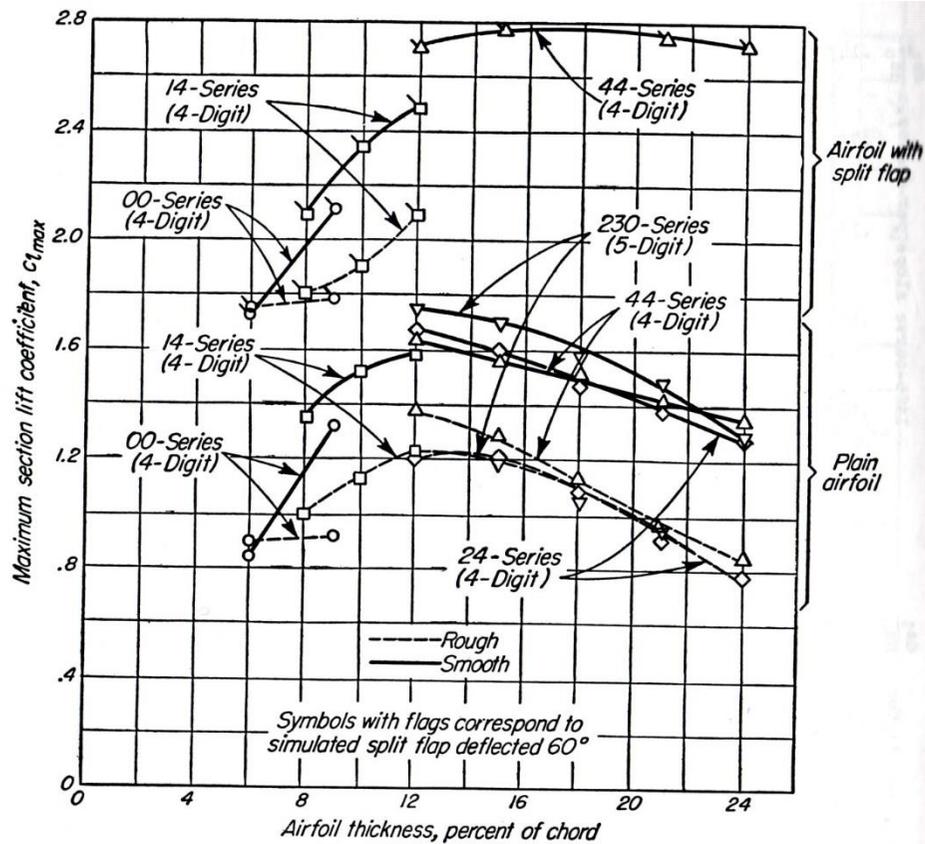


(c) NACA 64- series.

Fig. 57. Variation of lift-curve slope with airfoil thickness ratio and camber for a

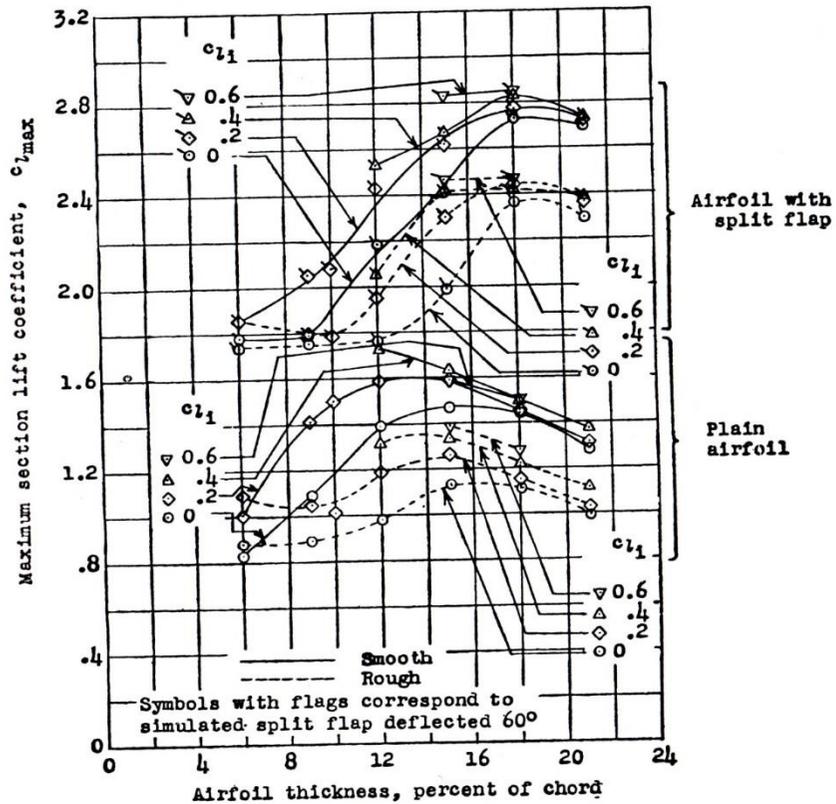


number of NACA airfoil sections in both the smooth and rough conditions.  $R, 6 \times 10^6$ .



(a) NACA four- and five-digit series.

FIG. 58. Variation of maximum section lift coefficient with airfoil thickness ratio and camber for several NACA airfoil sections with and without simulated split flaps and standard roughness.  $R, 6 \times 10^6$ .



(b) NACA 63-series.

FIG. 58. (Continued)

sections that would be considered for practical application. This jog n

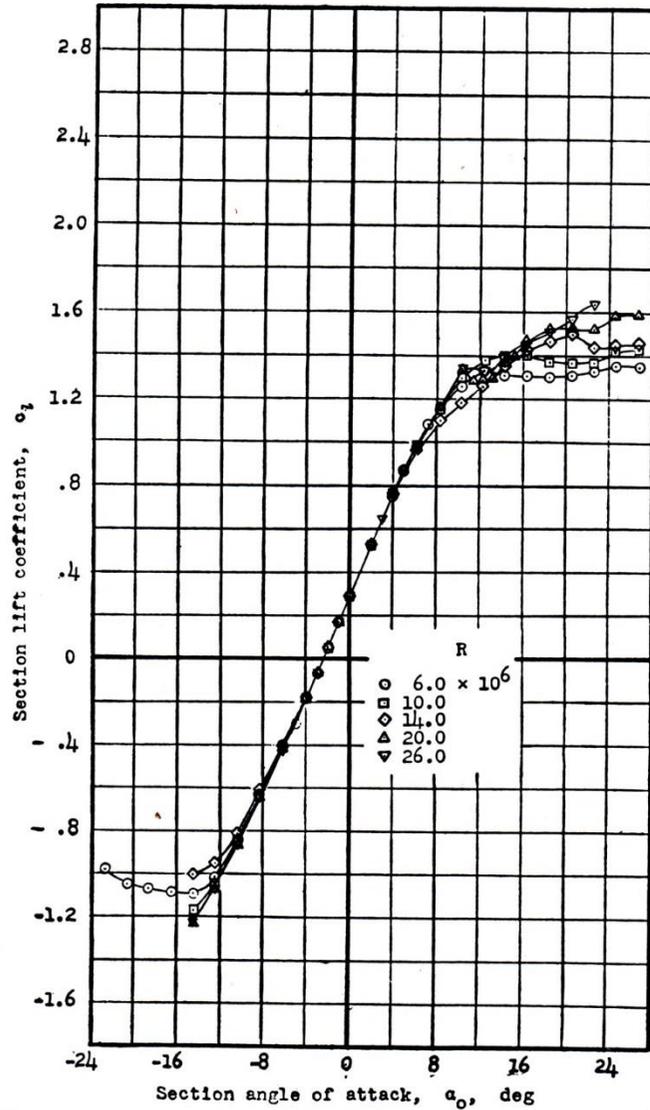


FIG. 61. Lift and drag characteristics of the NACA 63(420)-422 airfoil at high Reynolds number; TDT tests 228 and 255.

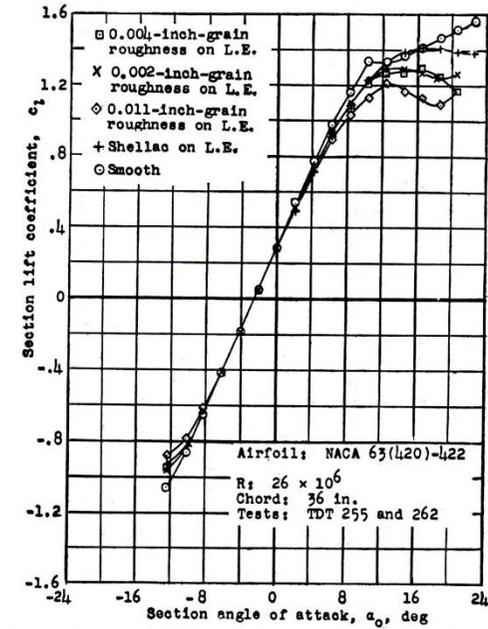


FIG. 62. Lift characteristics of a NACA 63(420)-422 airfoil with various degrees of roughness at the leading edge.

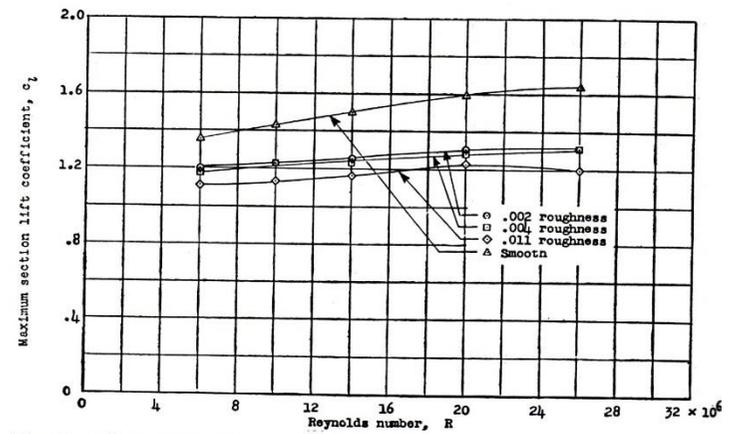


FIG. 63. Effects of Reynolds number on maximum section lift coefficient  $c_l$  of the NACA 63(420)-422 airfoil with roughened and smooth leading edge.

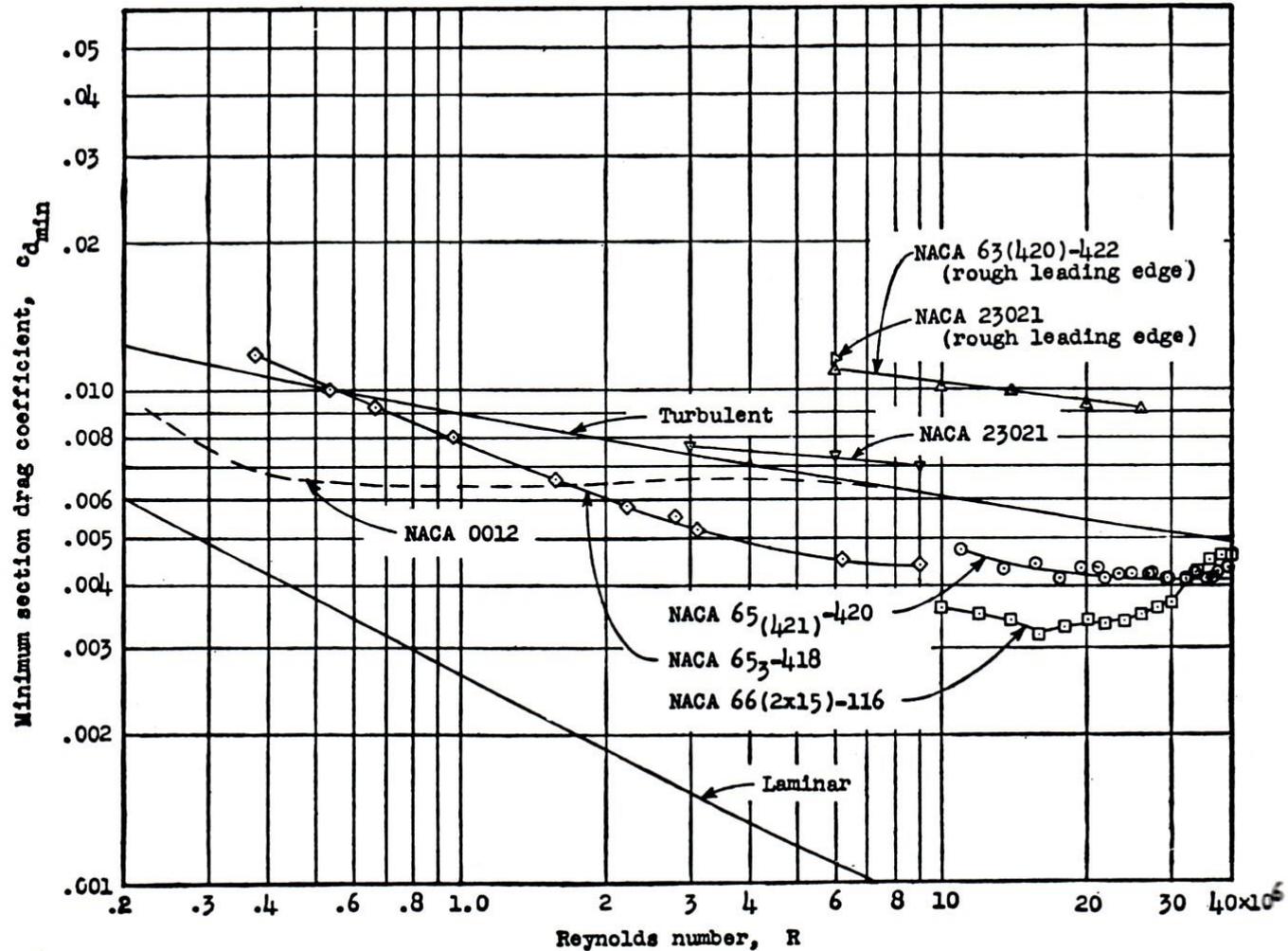


FIG. 66. Variation of minimum drag coefficient with Reynolds number for several airfoils, together with laminar and turbulent skin-friction coefficients for a flat plate.

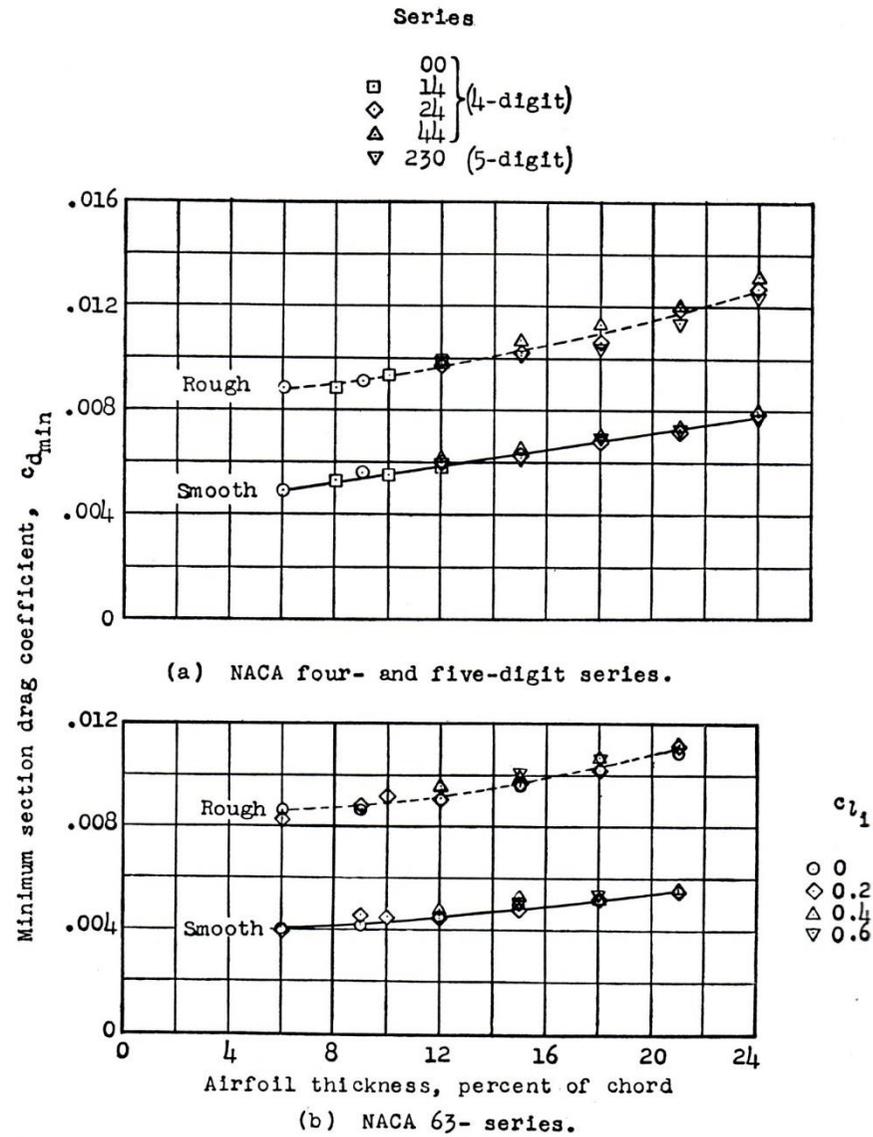


FIG. 68. Variation of section minimum drag coefficient with airfoil thickness ratio for several NACA airfoil sections of different cambers in both smooth and rough conditions.  $R, 6 \times 10^6$ .

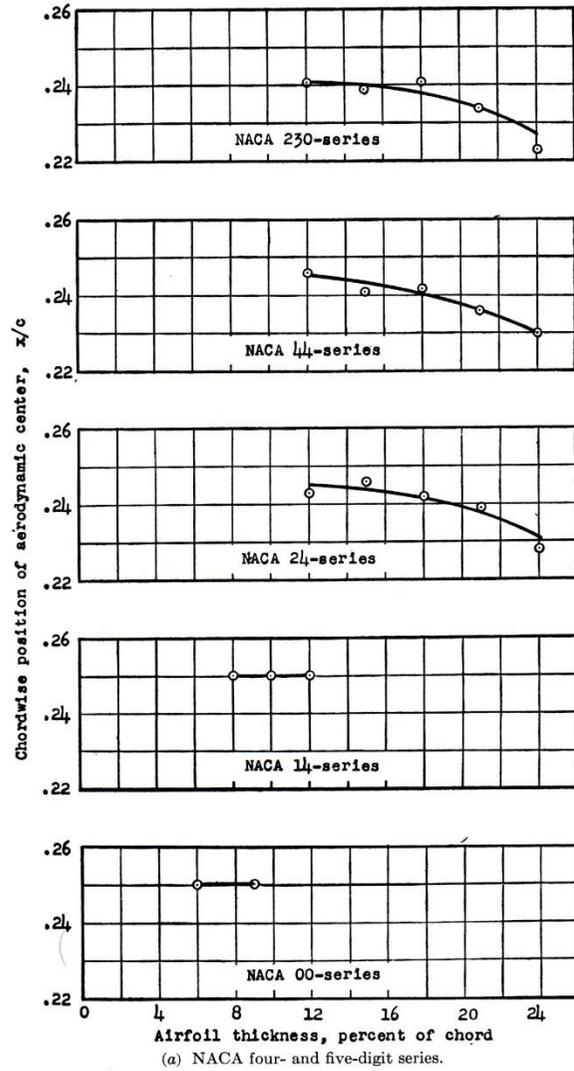


Fig. 94. Variation of section chordwise position of the aerodynamic center with airfoil thickness ratio for several NACA airfoil sections of different cambers.  $R, 6 \times 10^6$ .

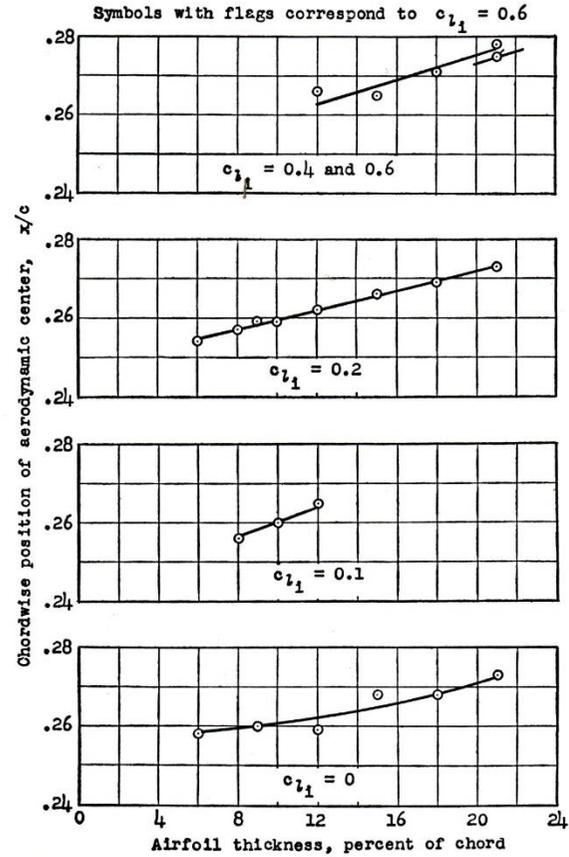
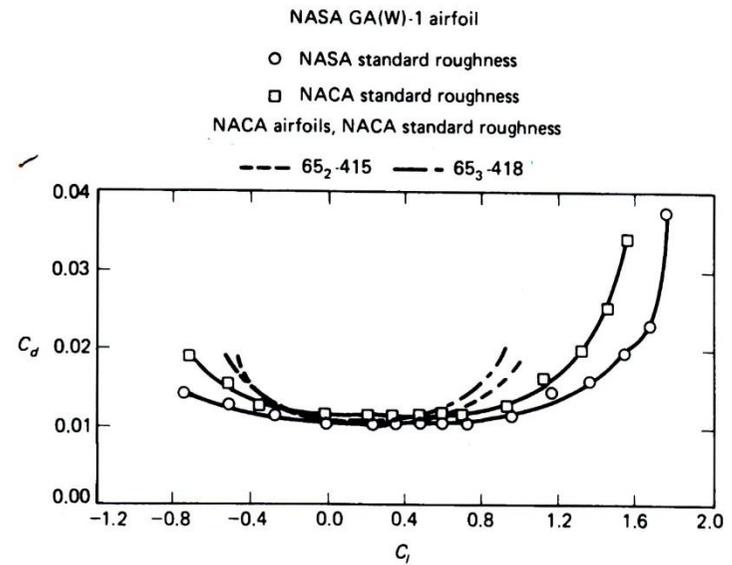
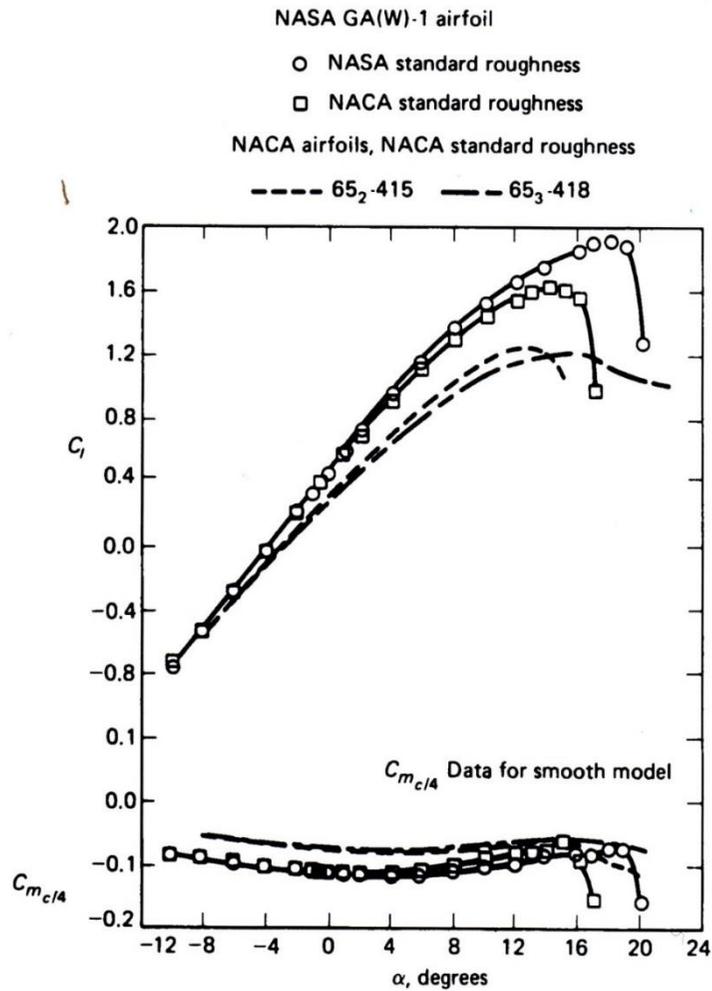


FIG. 94. (Continued)



R.J.McGhee, W.D.Beauley,- Low Speed Aerodynamic Characteristics of a 17-Percent-Thick Section Designed for General Aviation Application (NASA TN D-7428, 1973) [from J.J.Bertin, M.L.Smith – Aerodynamics for Engineers, Printice-Hall International, Inc., 1998]